

An Unhelpful Introduction to Electricity & Magnetism

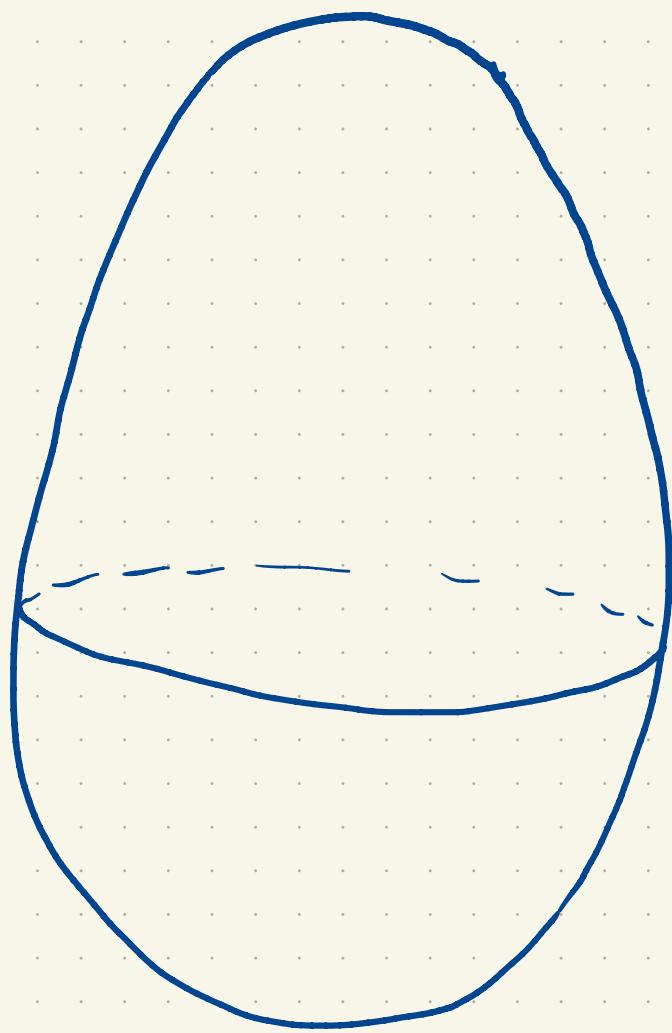
Part III^{1/2}: Eggs, Curvature and Two Maxwell Equations

Nov 24, 2020

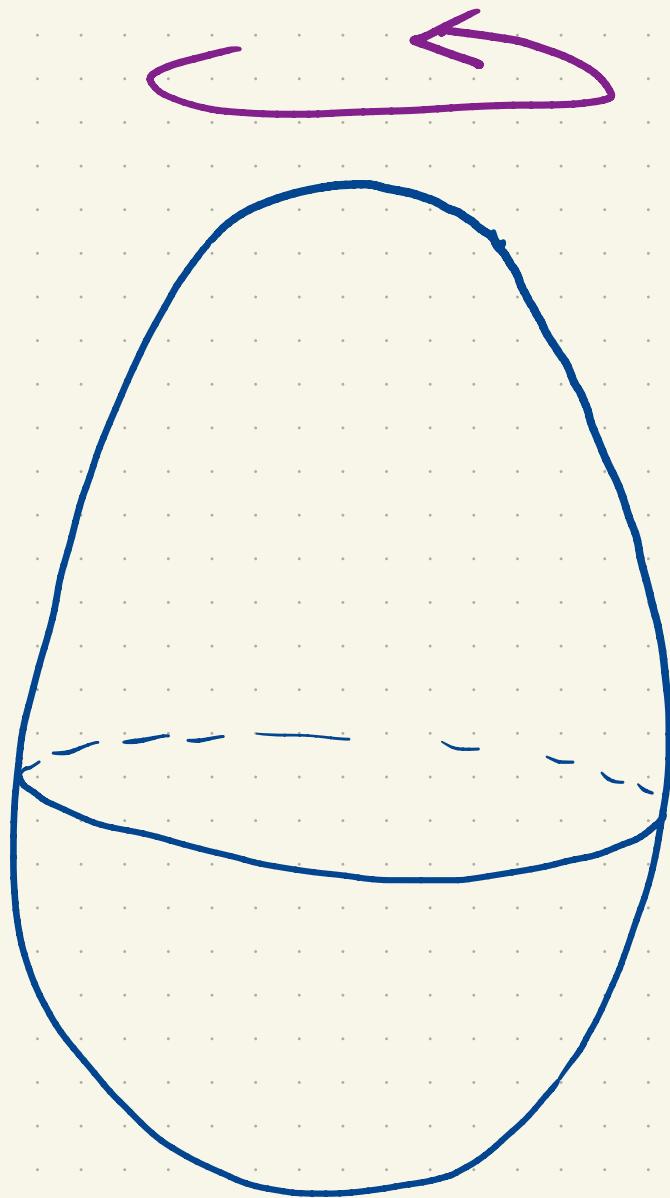
David Maxwell (damaxwell.github.io)

UAF Mathematics

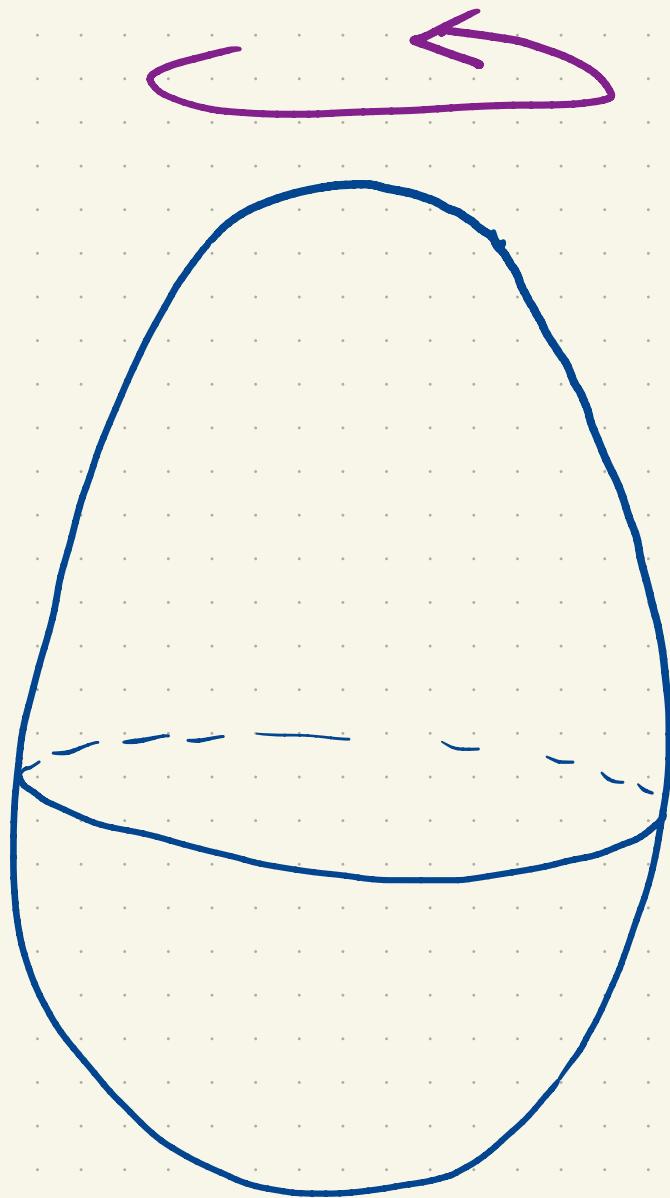
The Egg



The Egg



The Egg

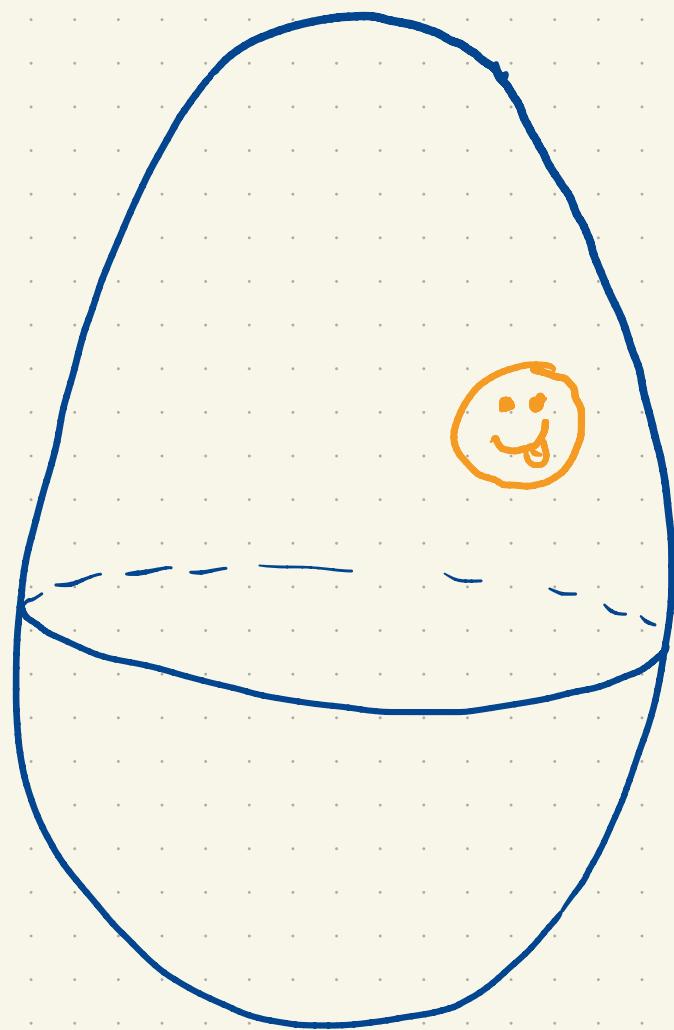
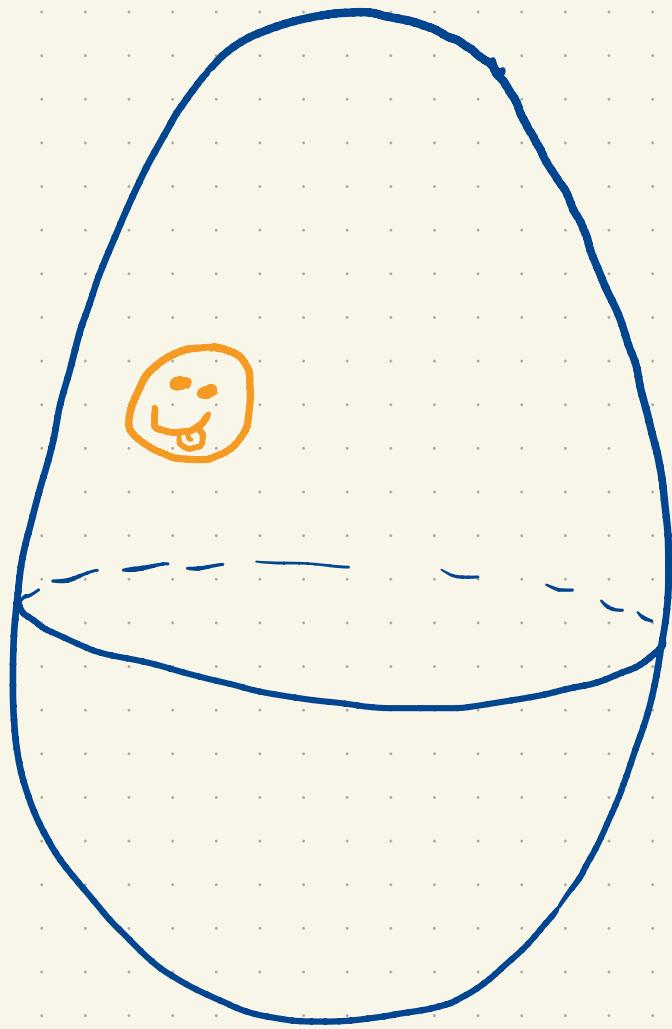


- Symmetry group:

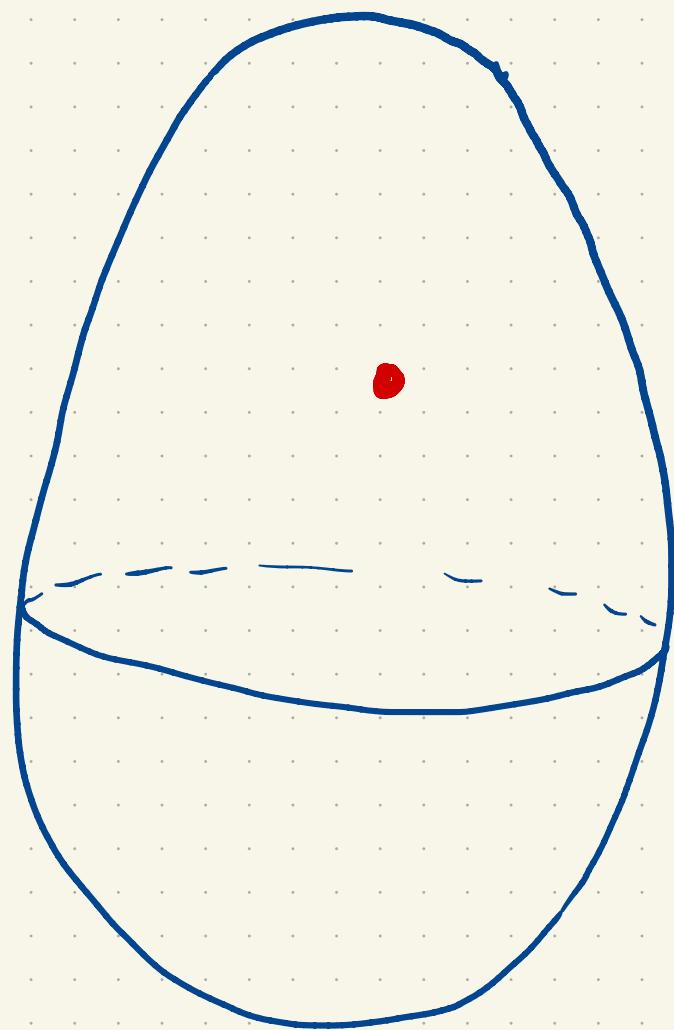
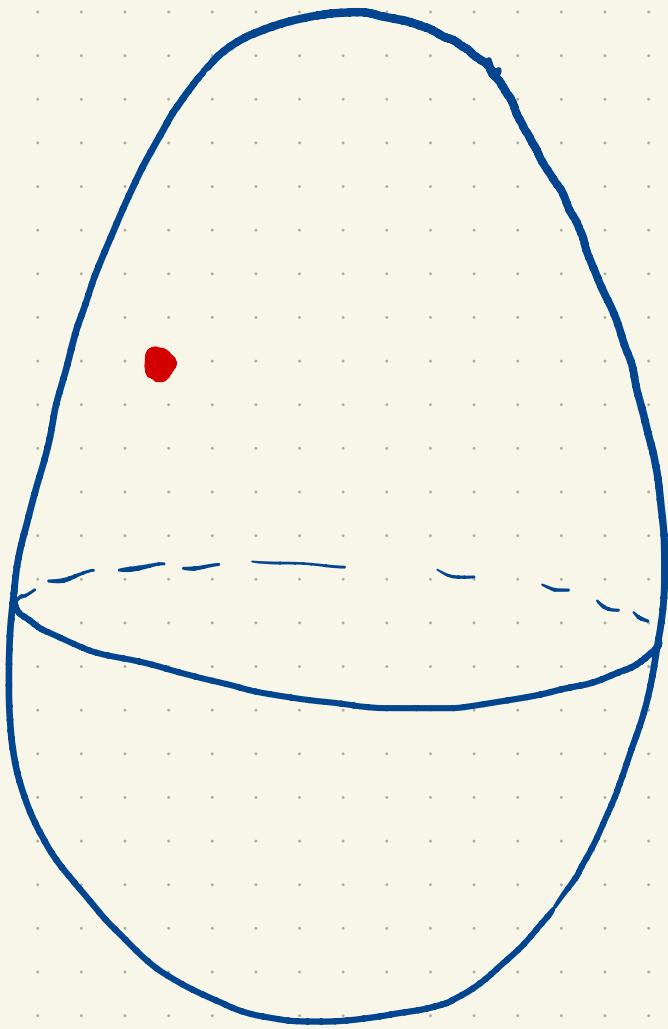
$$U(1)$$

$$\{e^{i\theta} : \theta \in \mathbb{R}\}$$

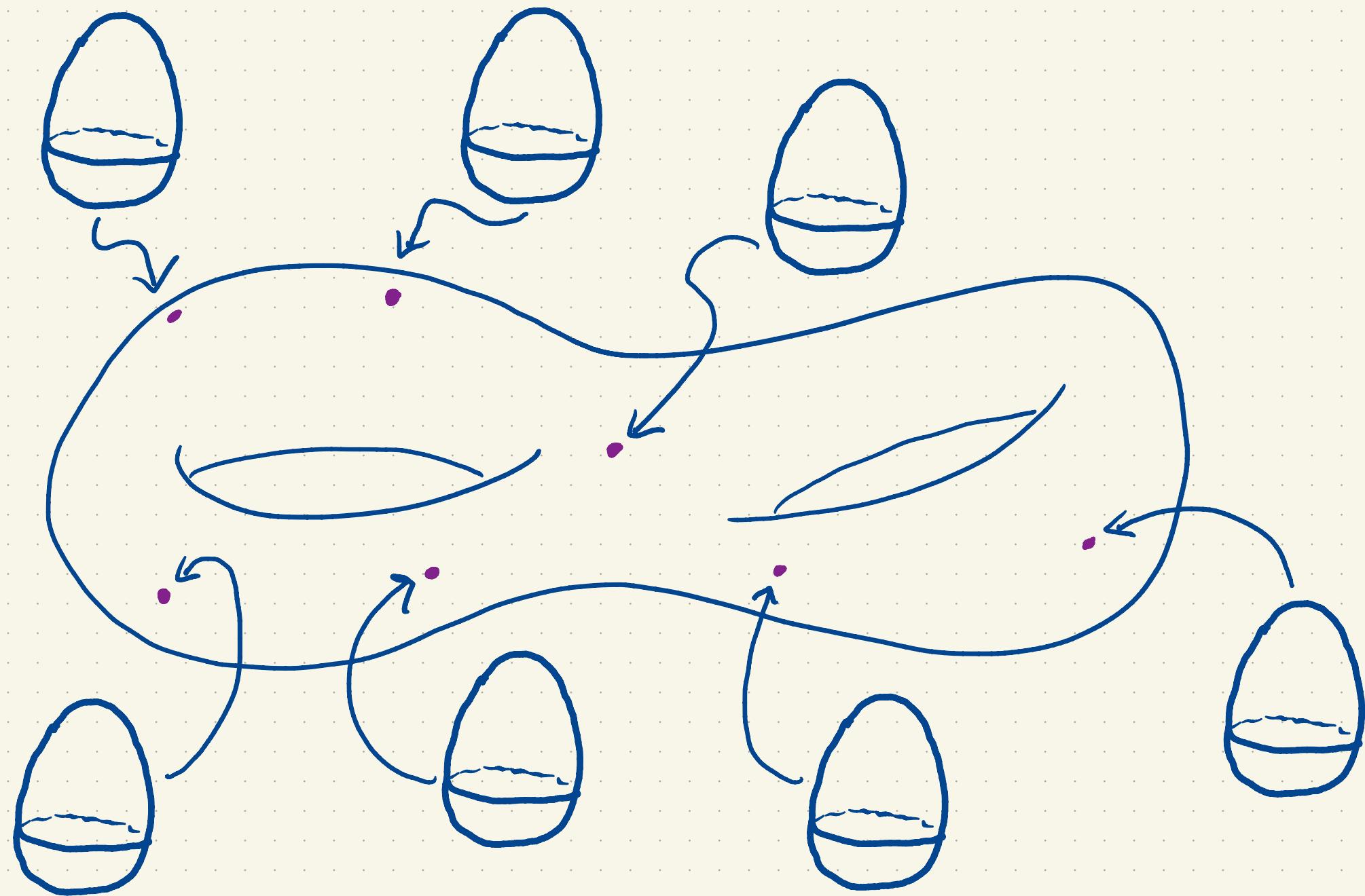
Congruence



Congruence

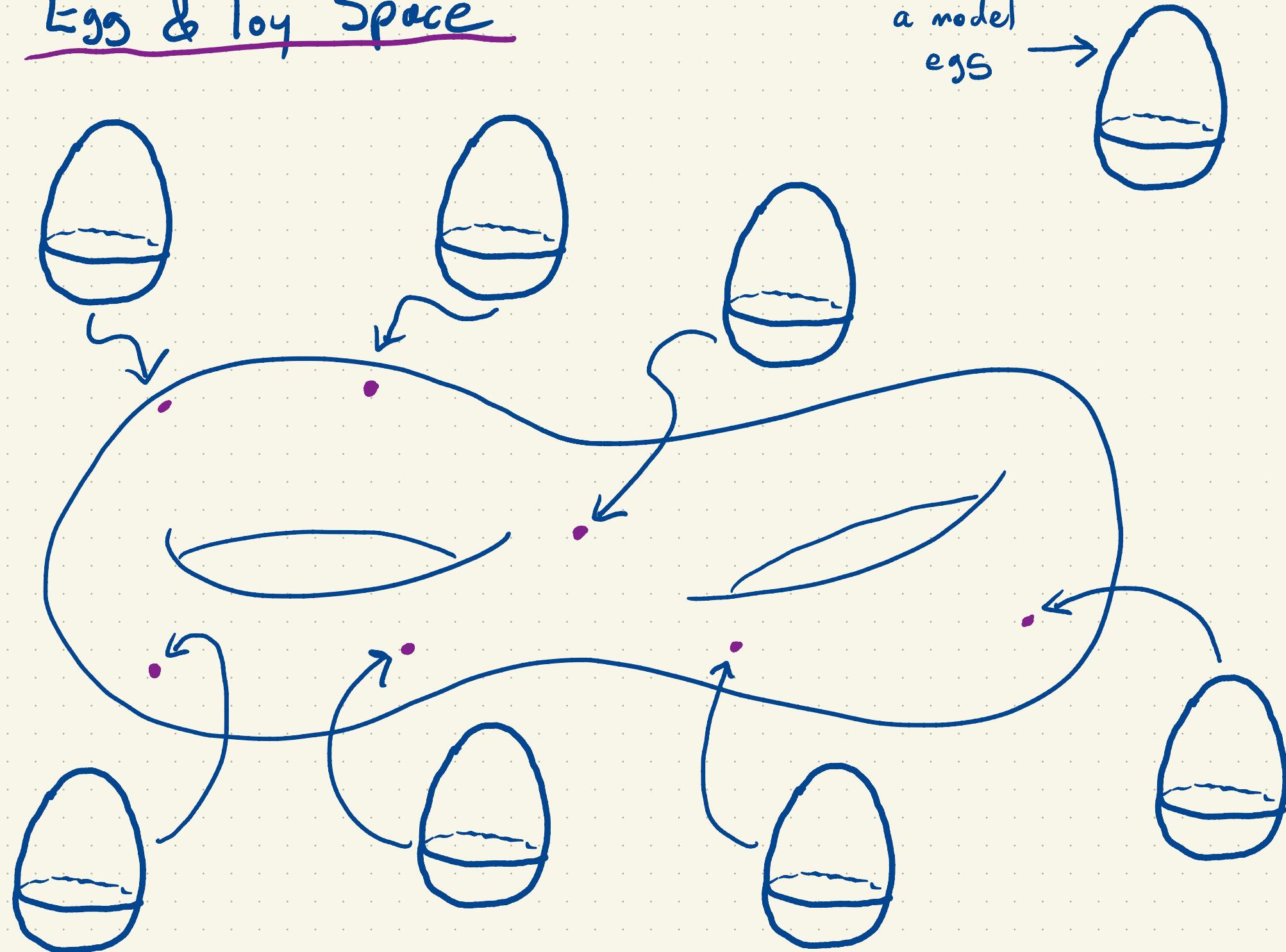


Eggs & Toy Space

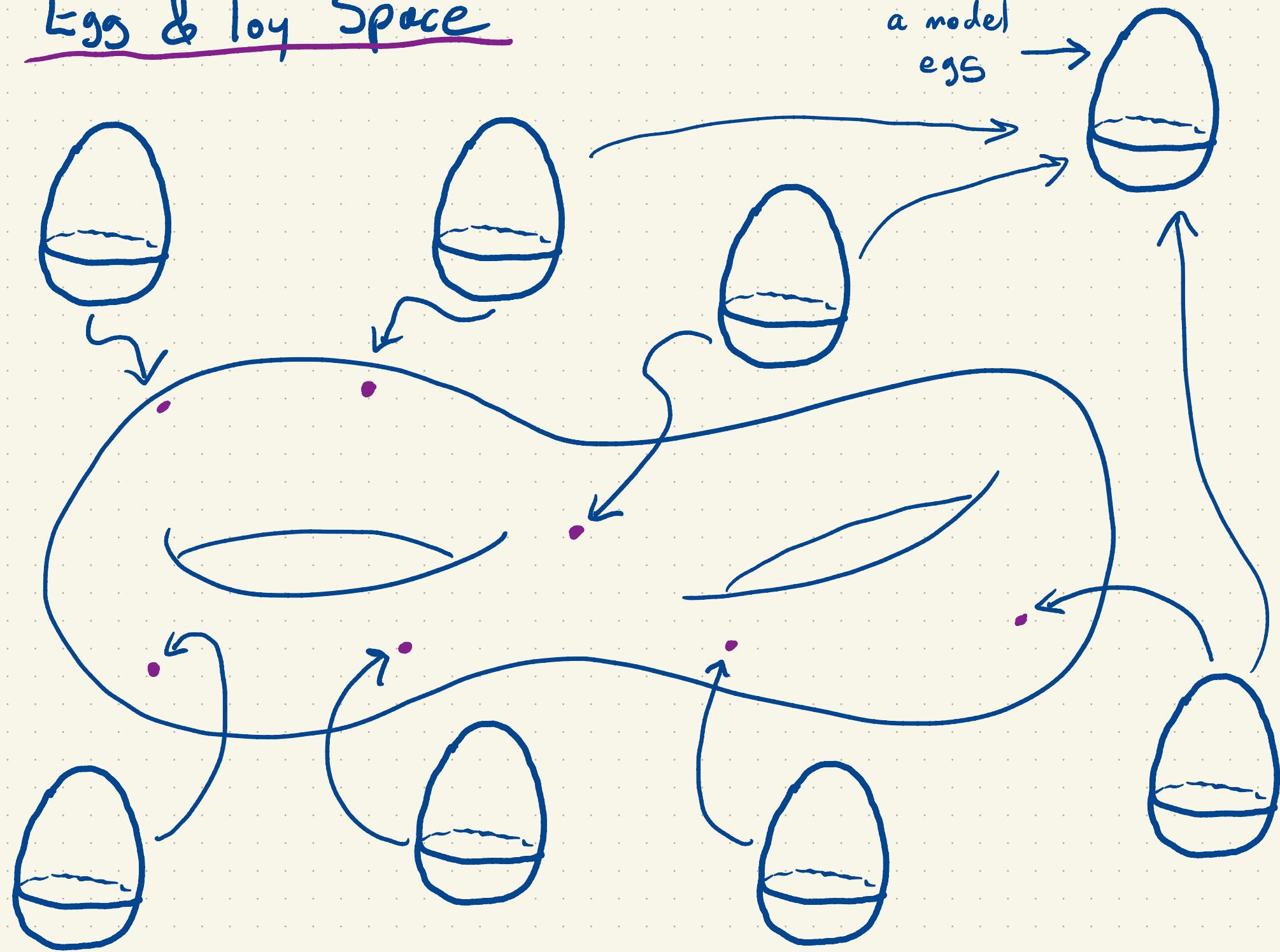


Eggs & Toy Space

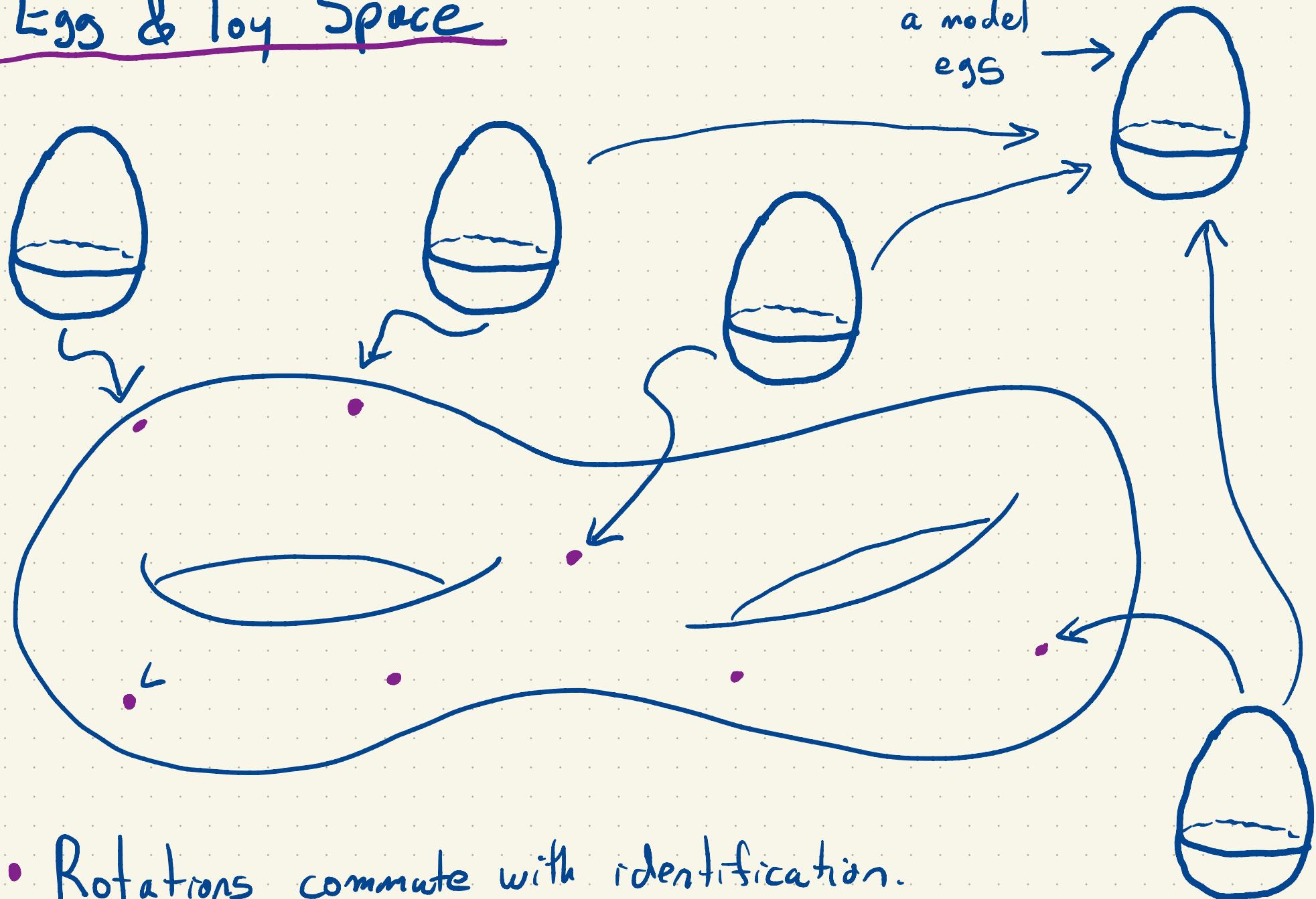
a model
egg



Eggs & Toy Space

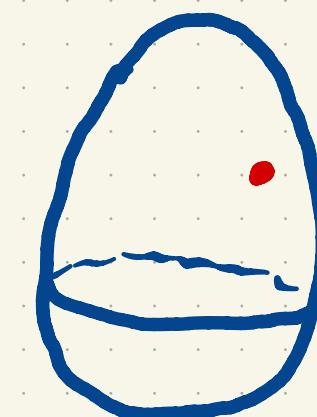
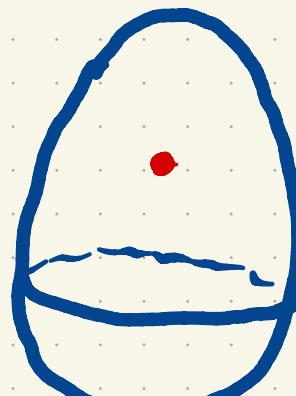
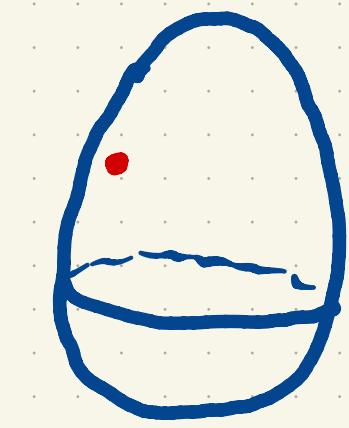
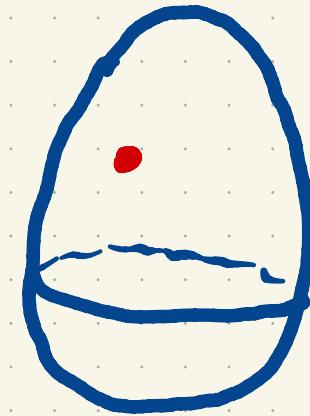


Eggs & Toy Space

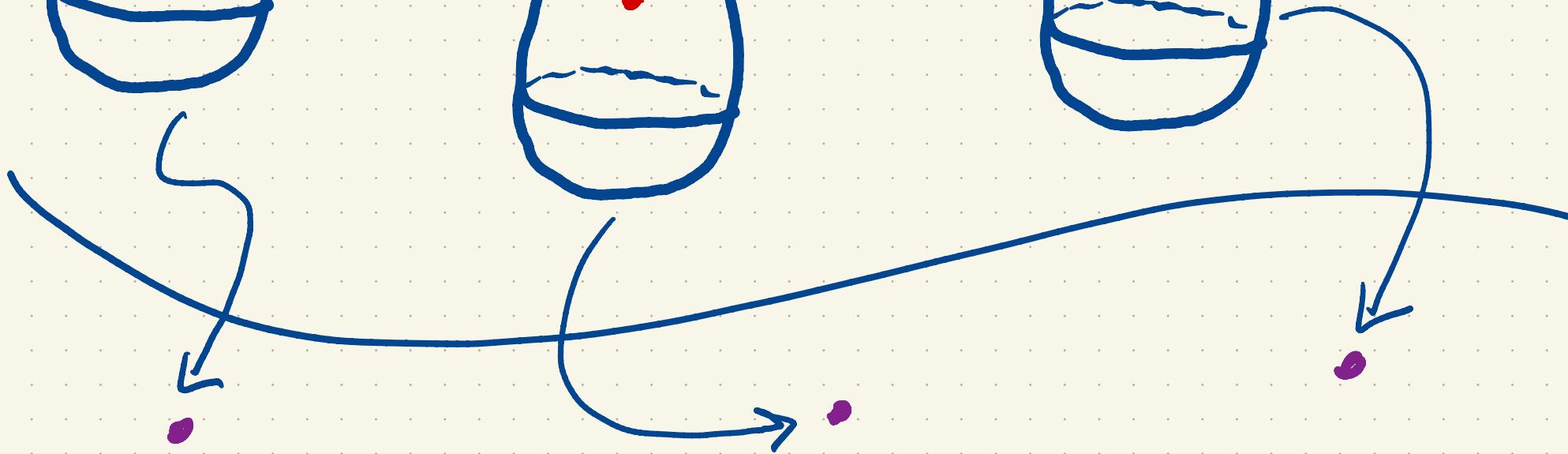
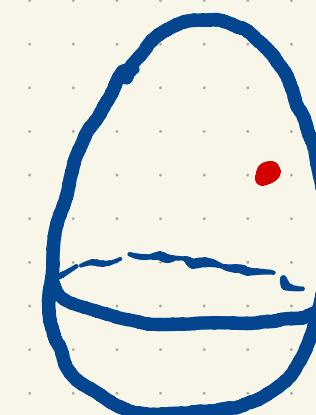
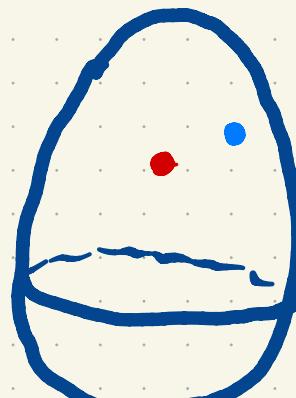
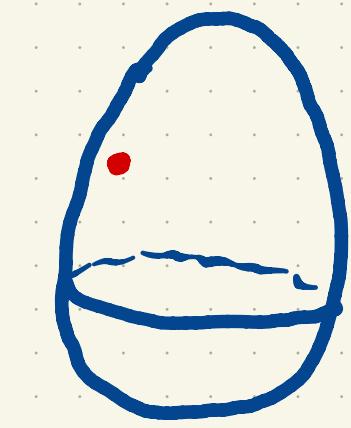
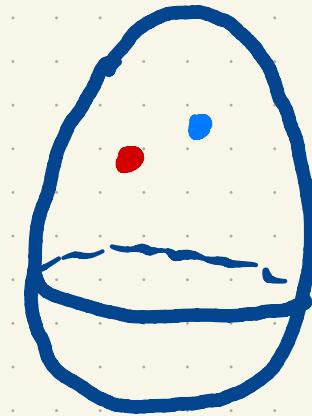


- Rotations commute with identification.
- Any two eggs have a circle's worth of identifications.

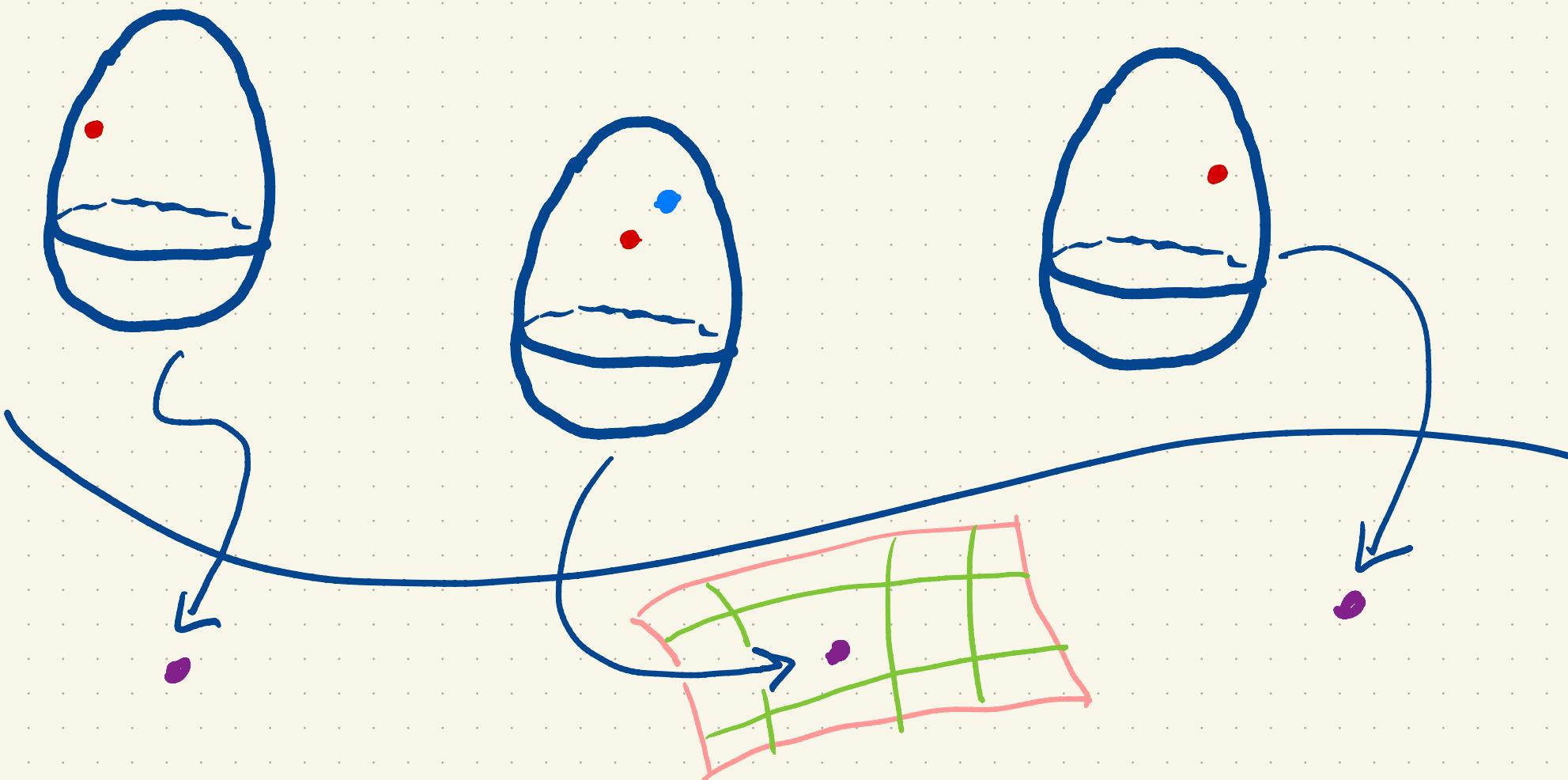
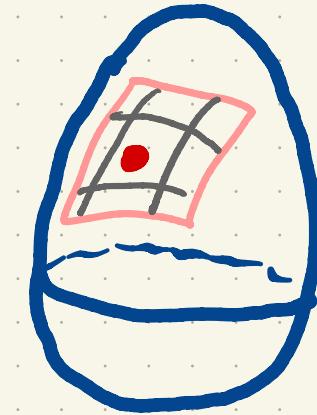
Egg Coordinates



Egg Coordinates

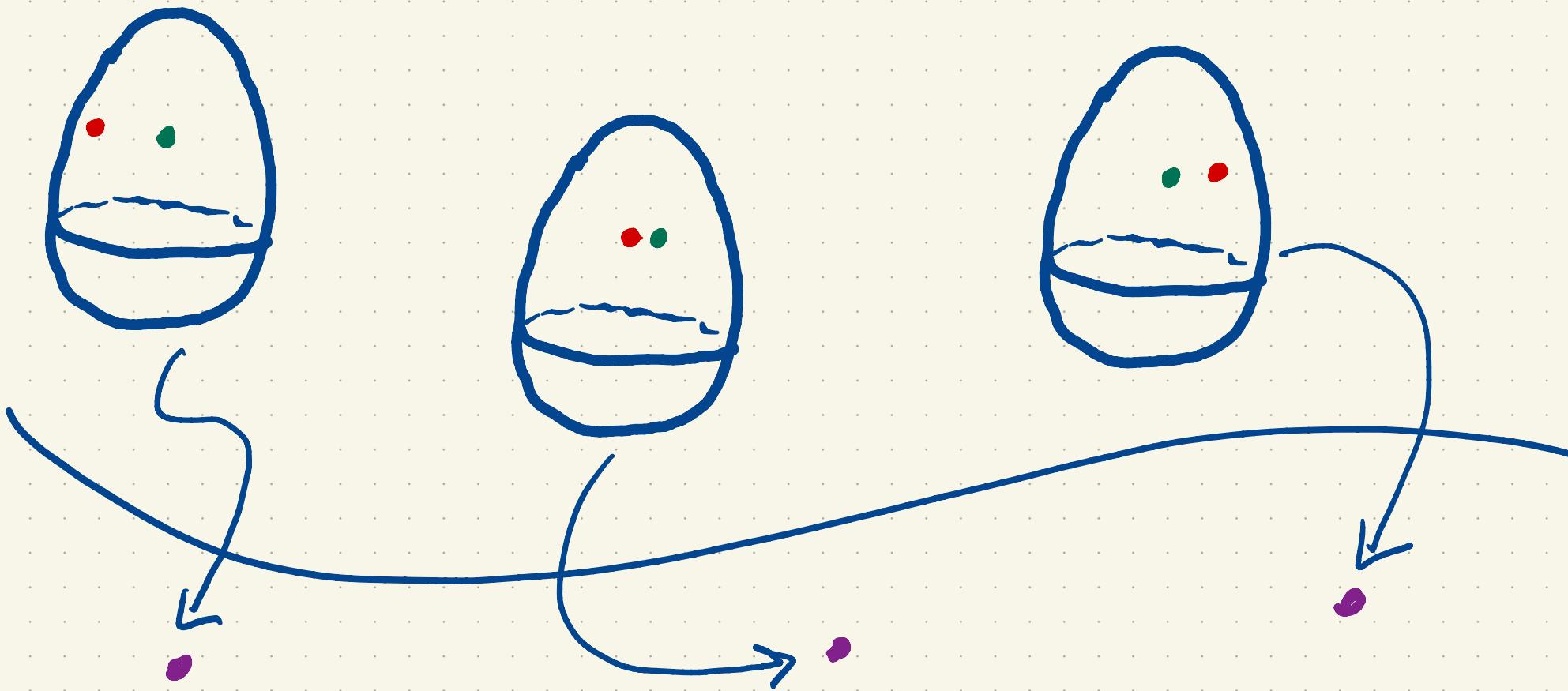
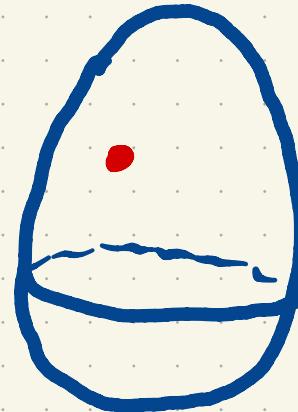


Egg Coordinates



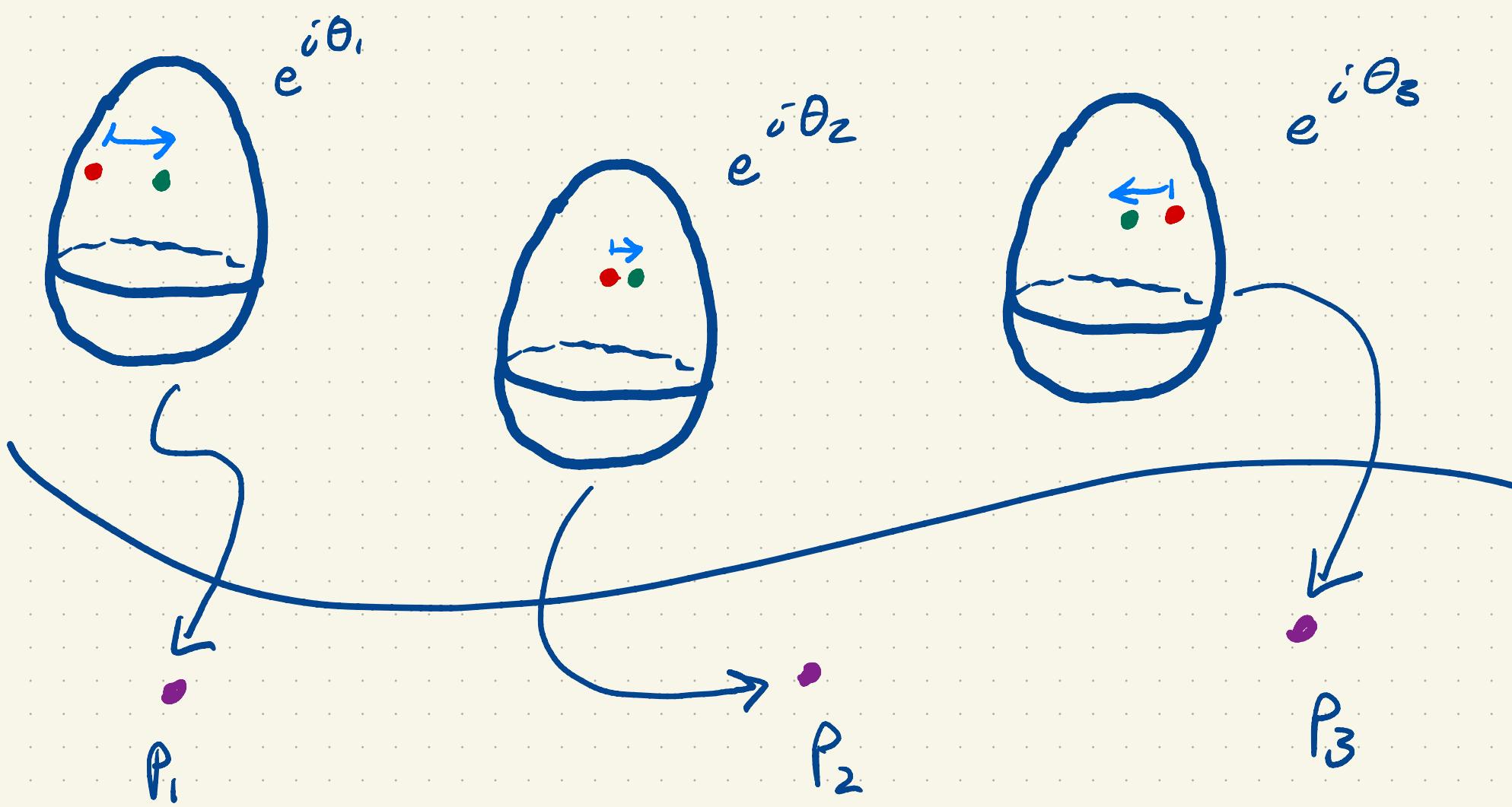
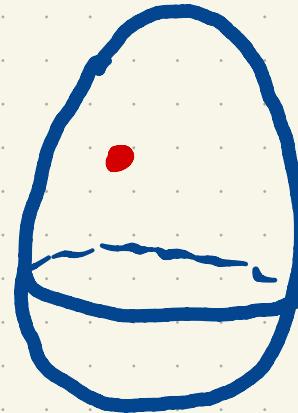
Egg Coordinates

- Two different egg coordinate systems.



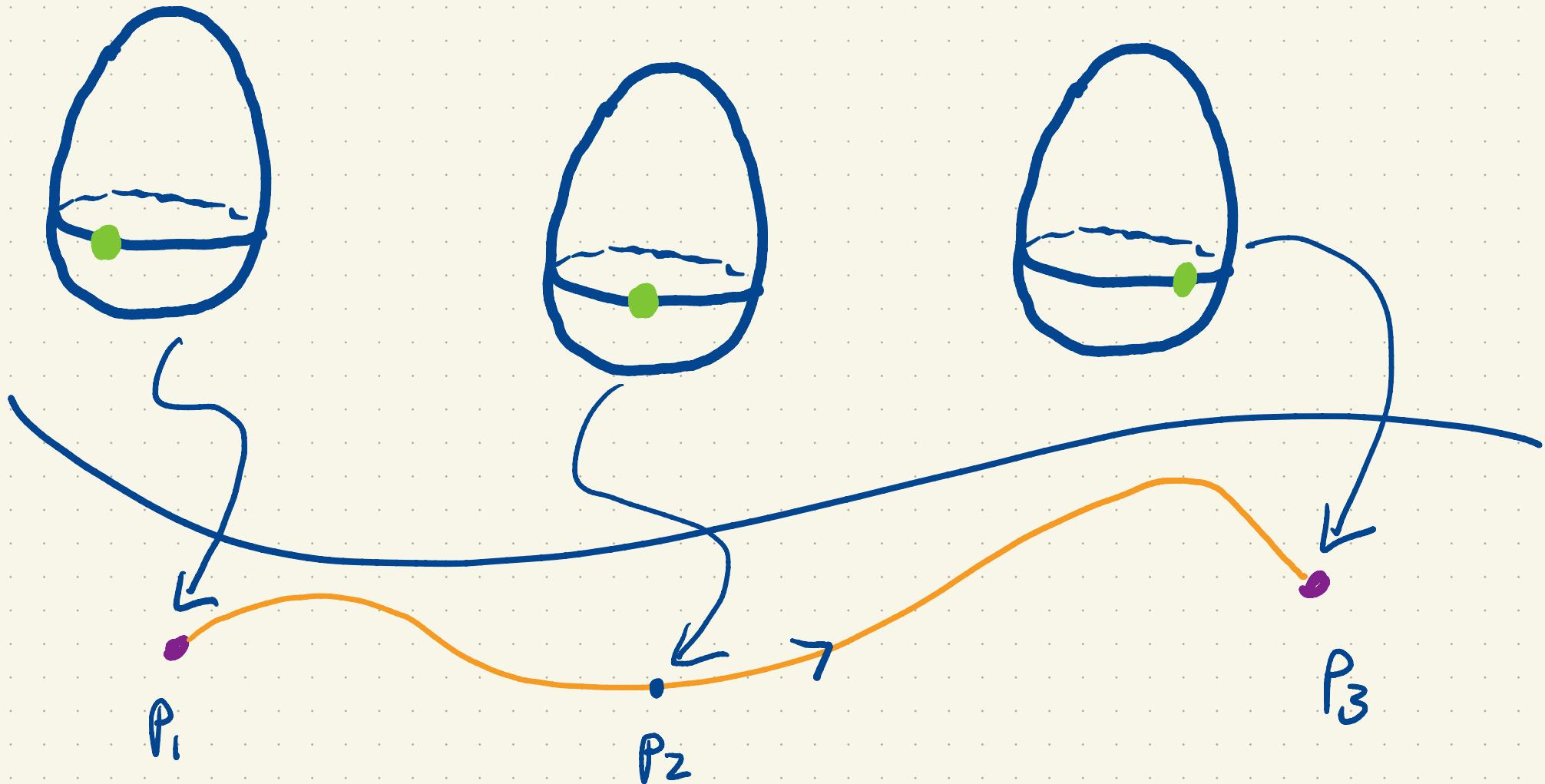
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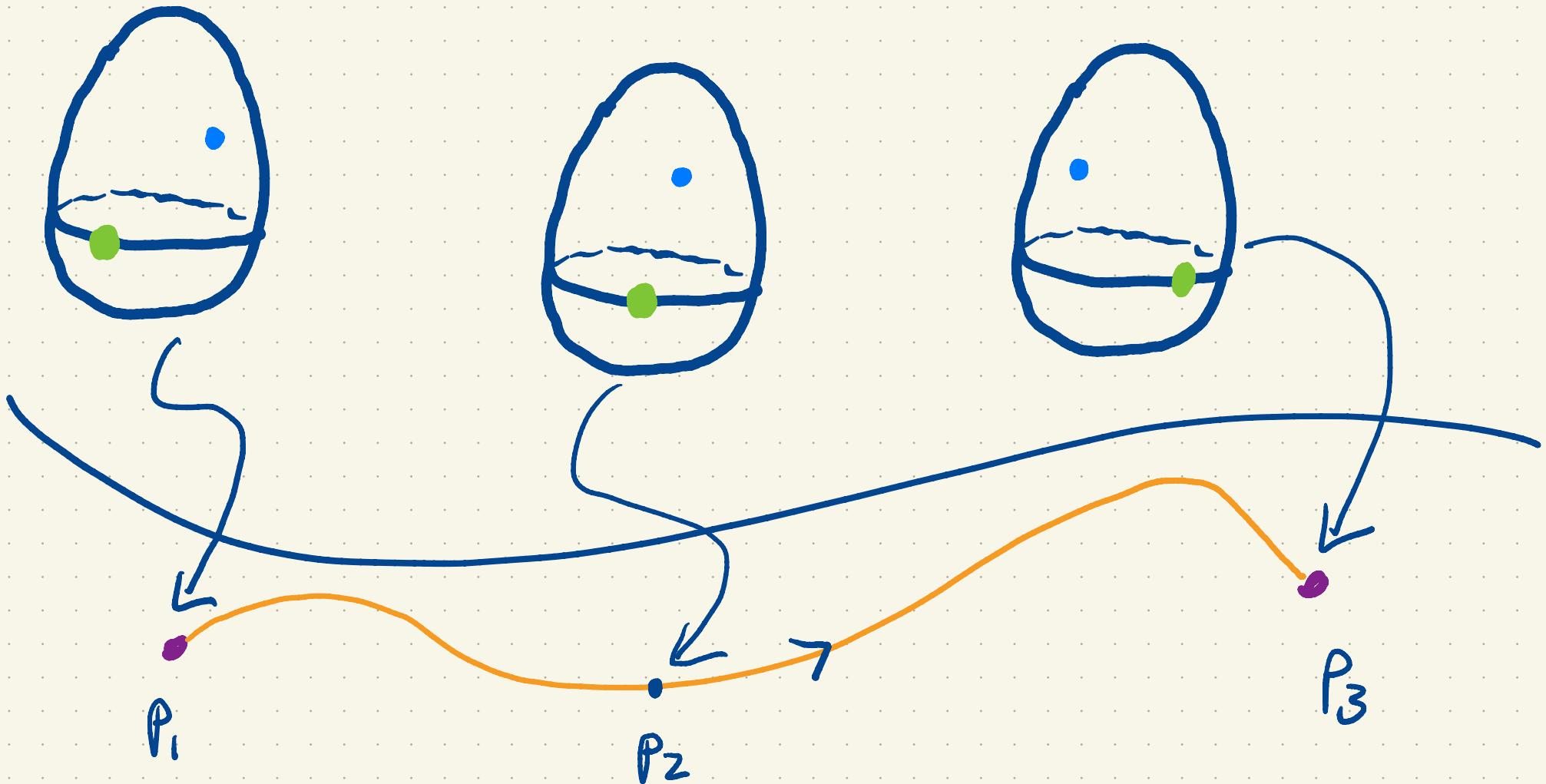
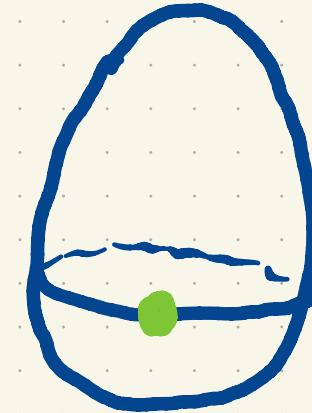
Egg Dragging

- Is an additional structure beyond the egg structure.



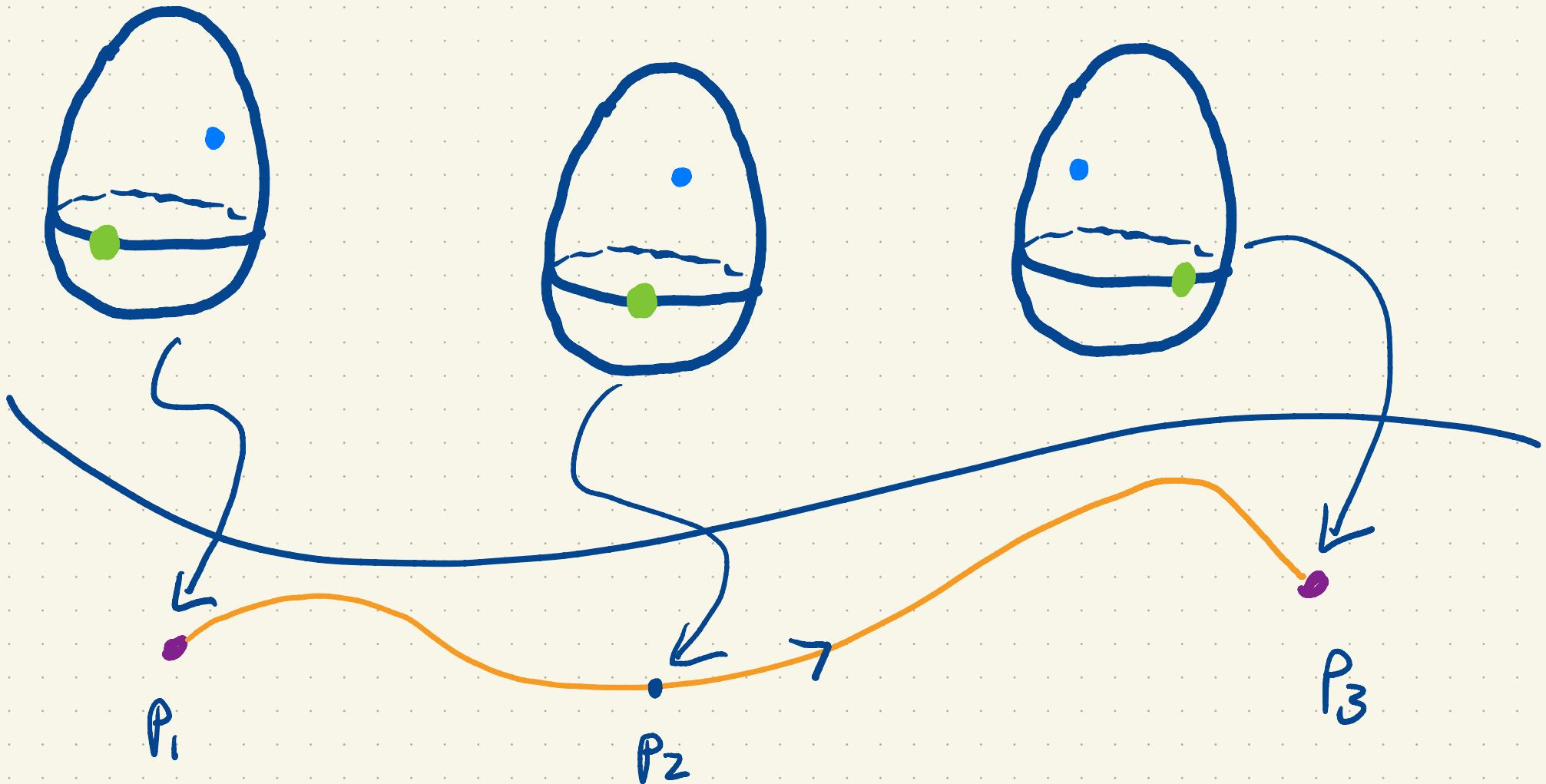
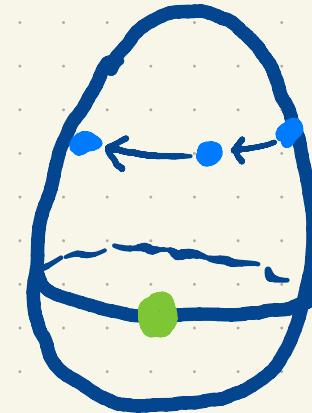
Egg Dragging

- It lets you measure change in the egg along a curve.



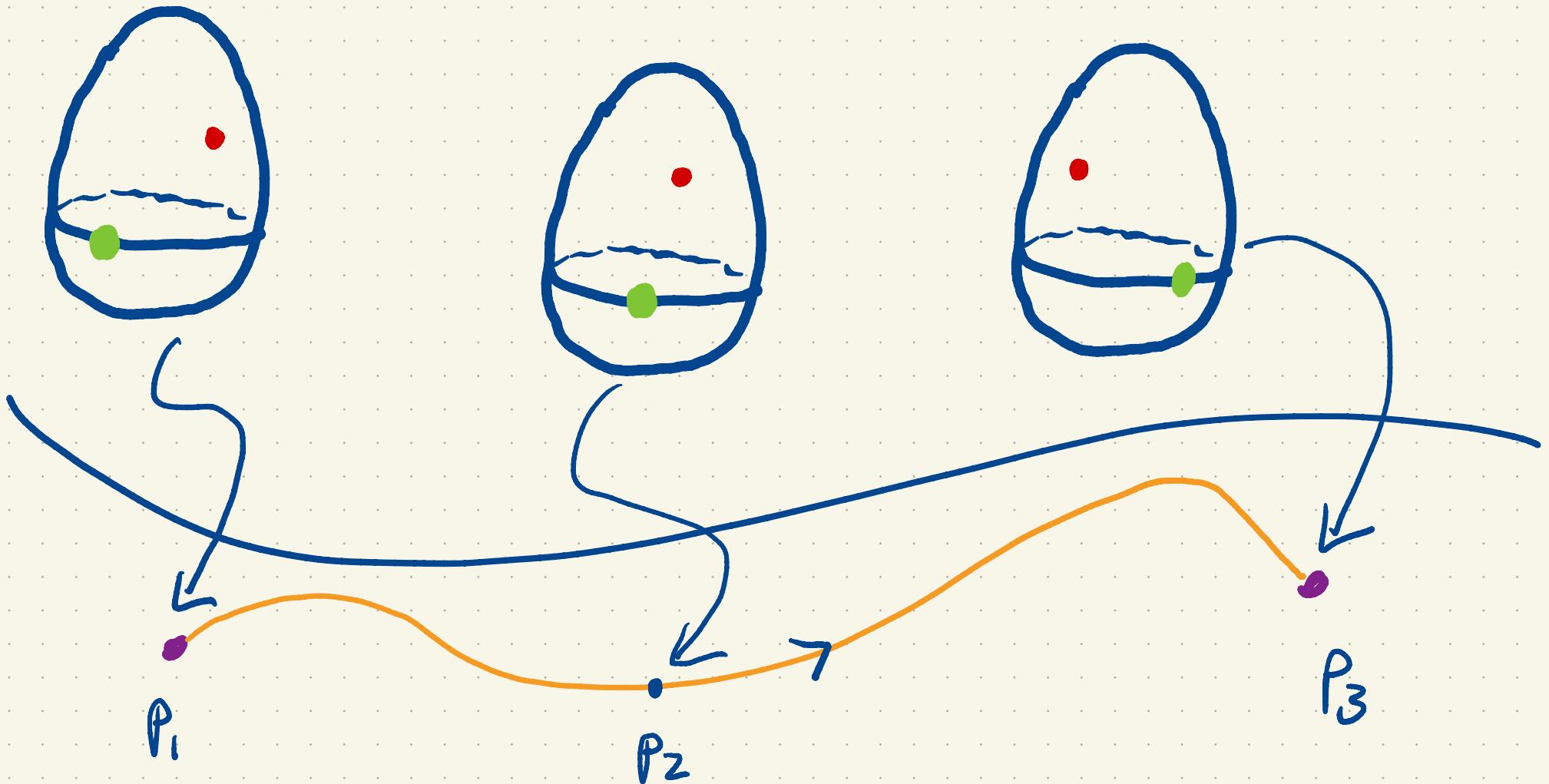
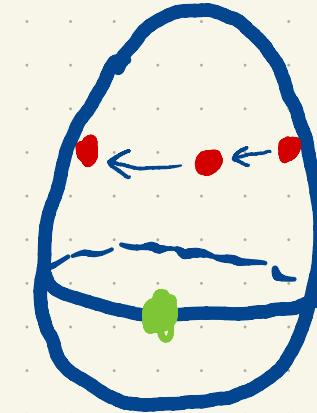
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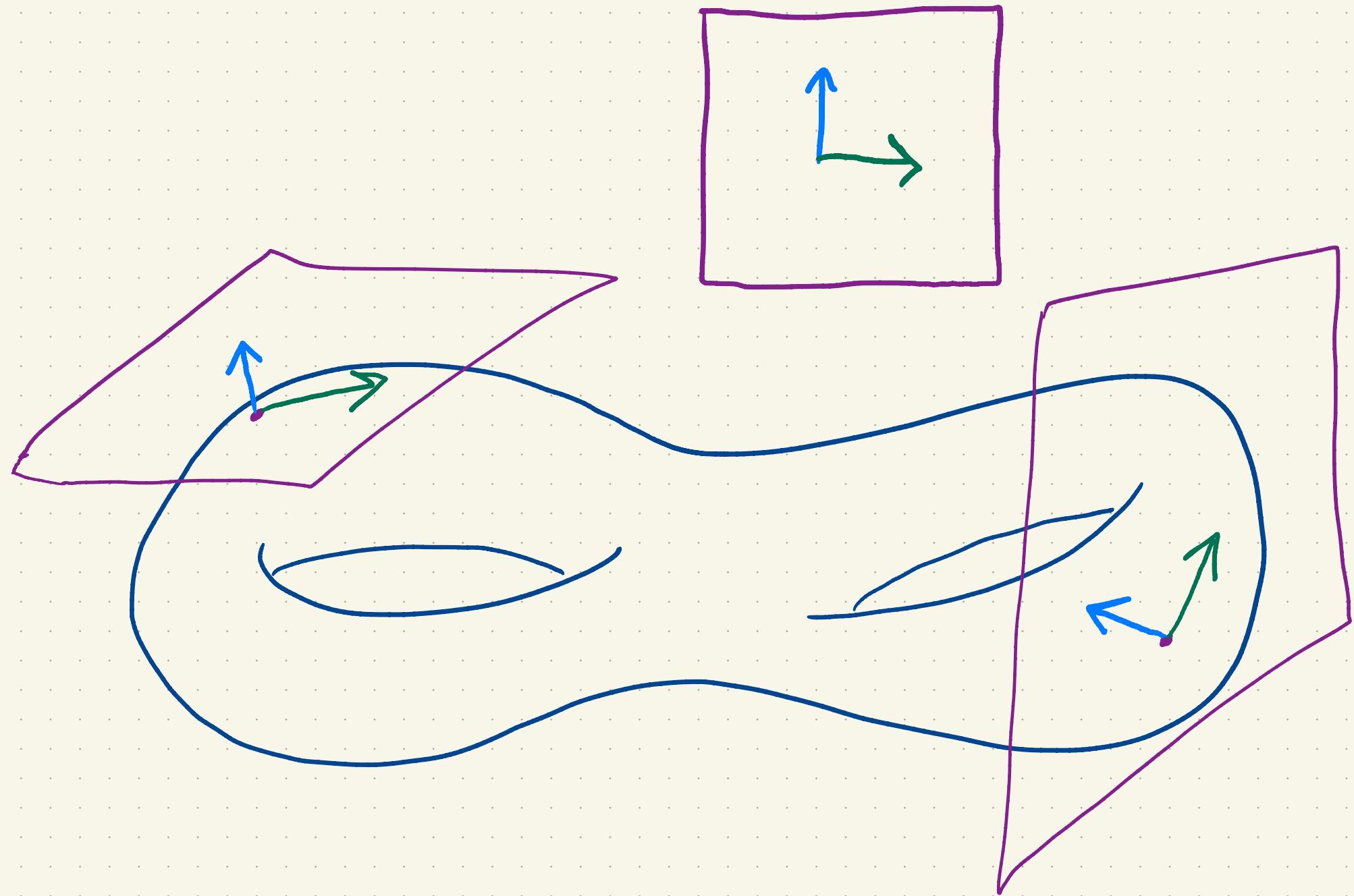


Egg Dragging

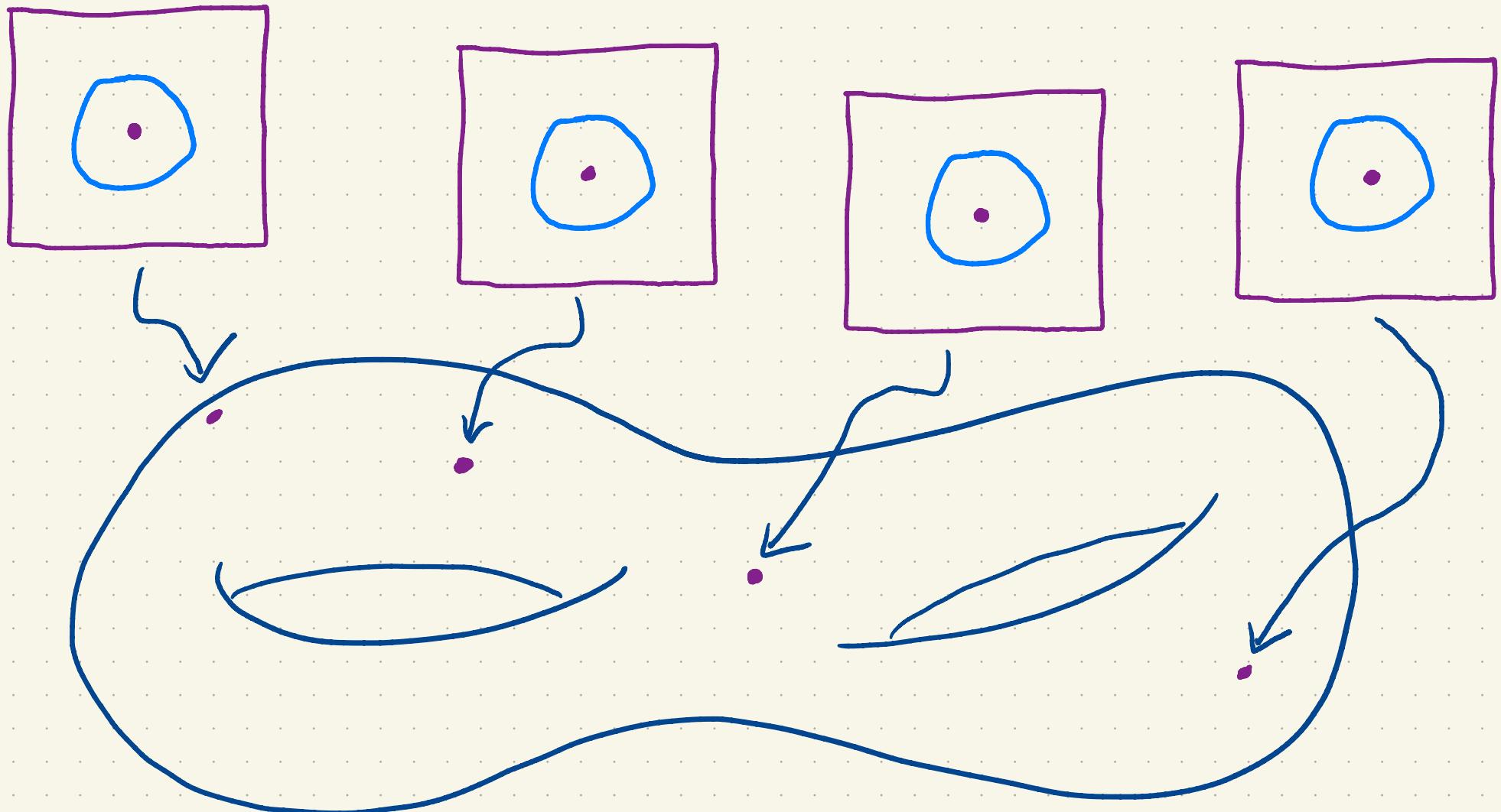
- If you pick egg coordinates it tells you "rate of rotation" of coordinates along the curve



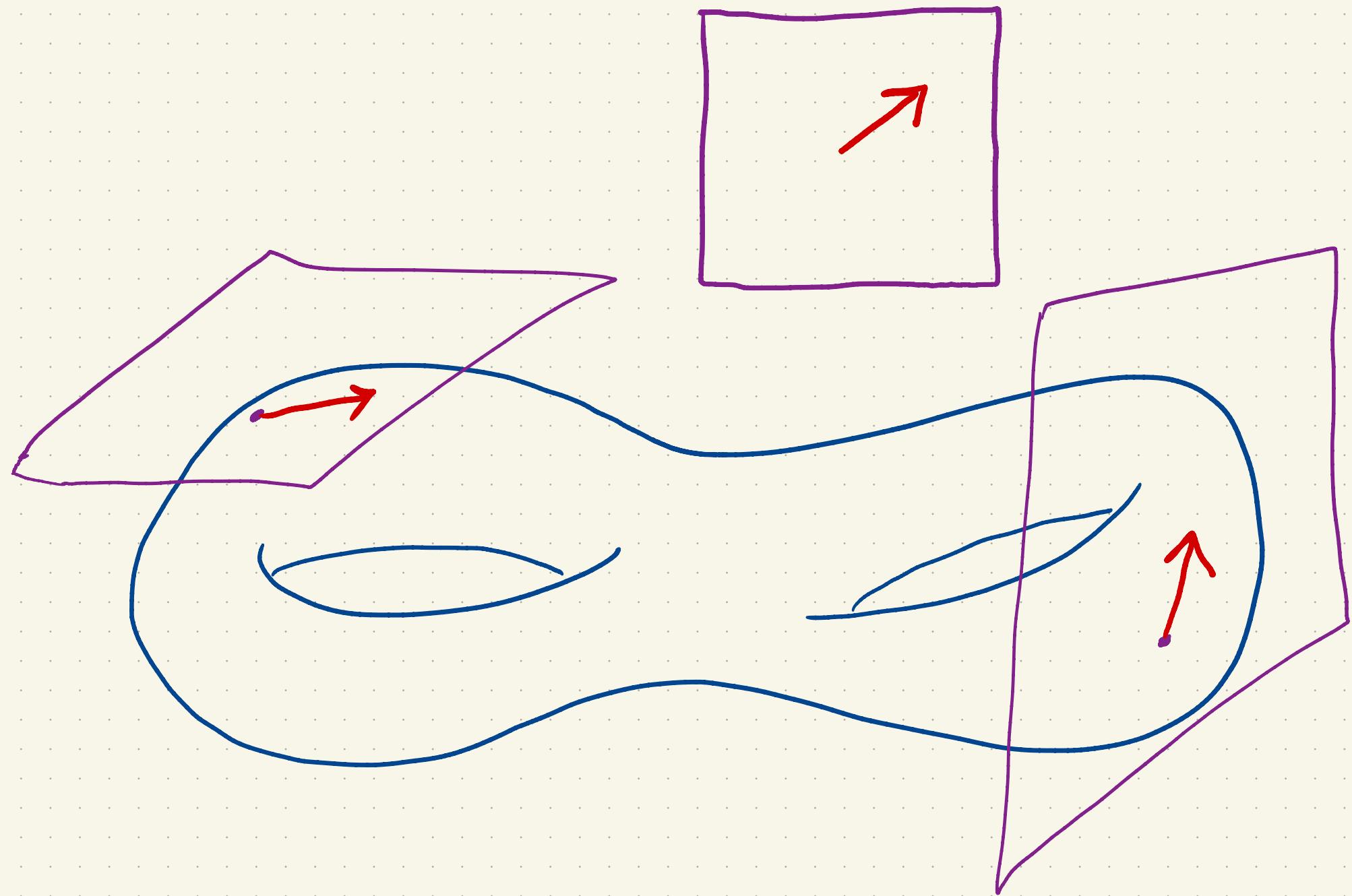
An Egg Can be a Plane



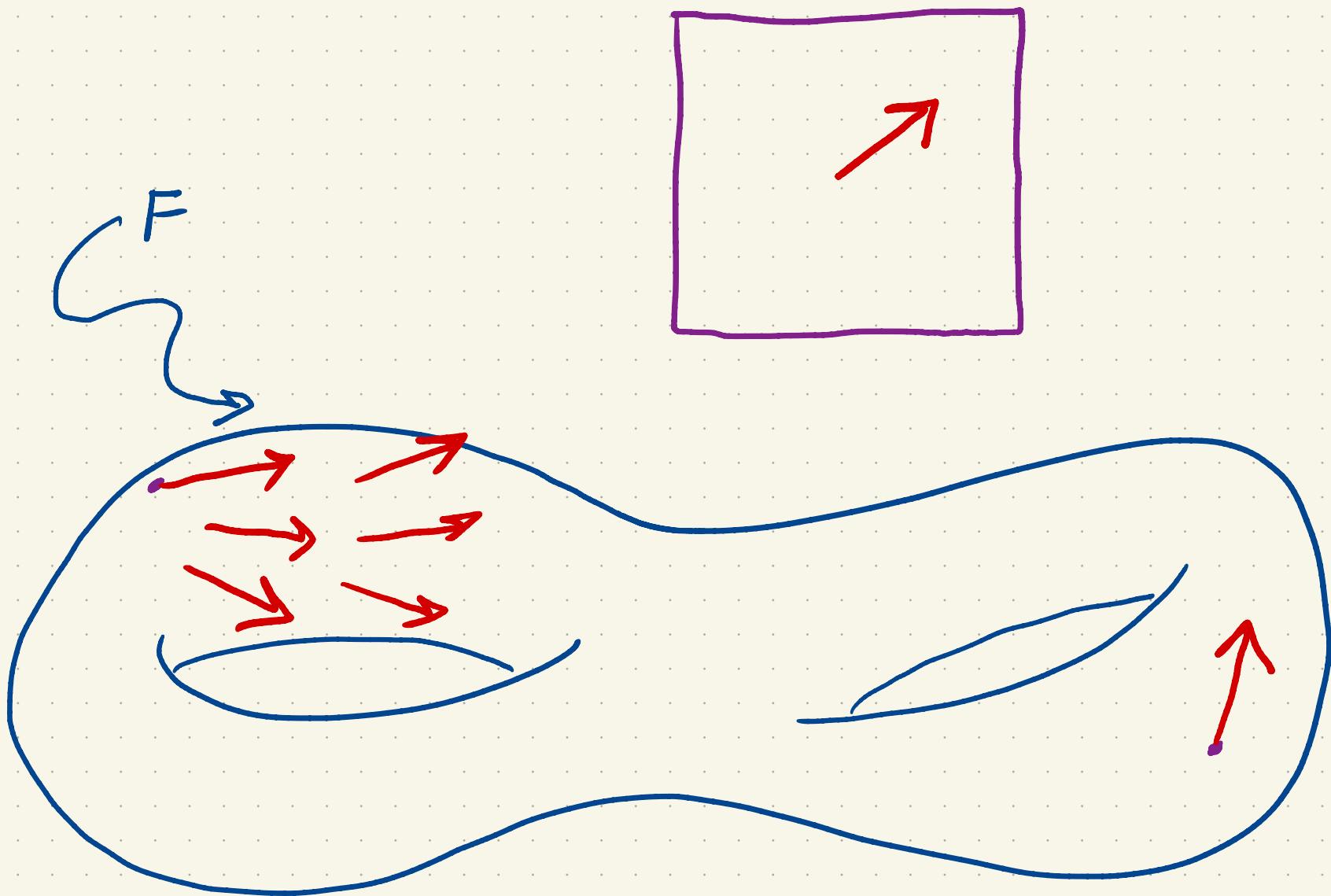
An Egg Can be a Plane



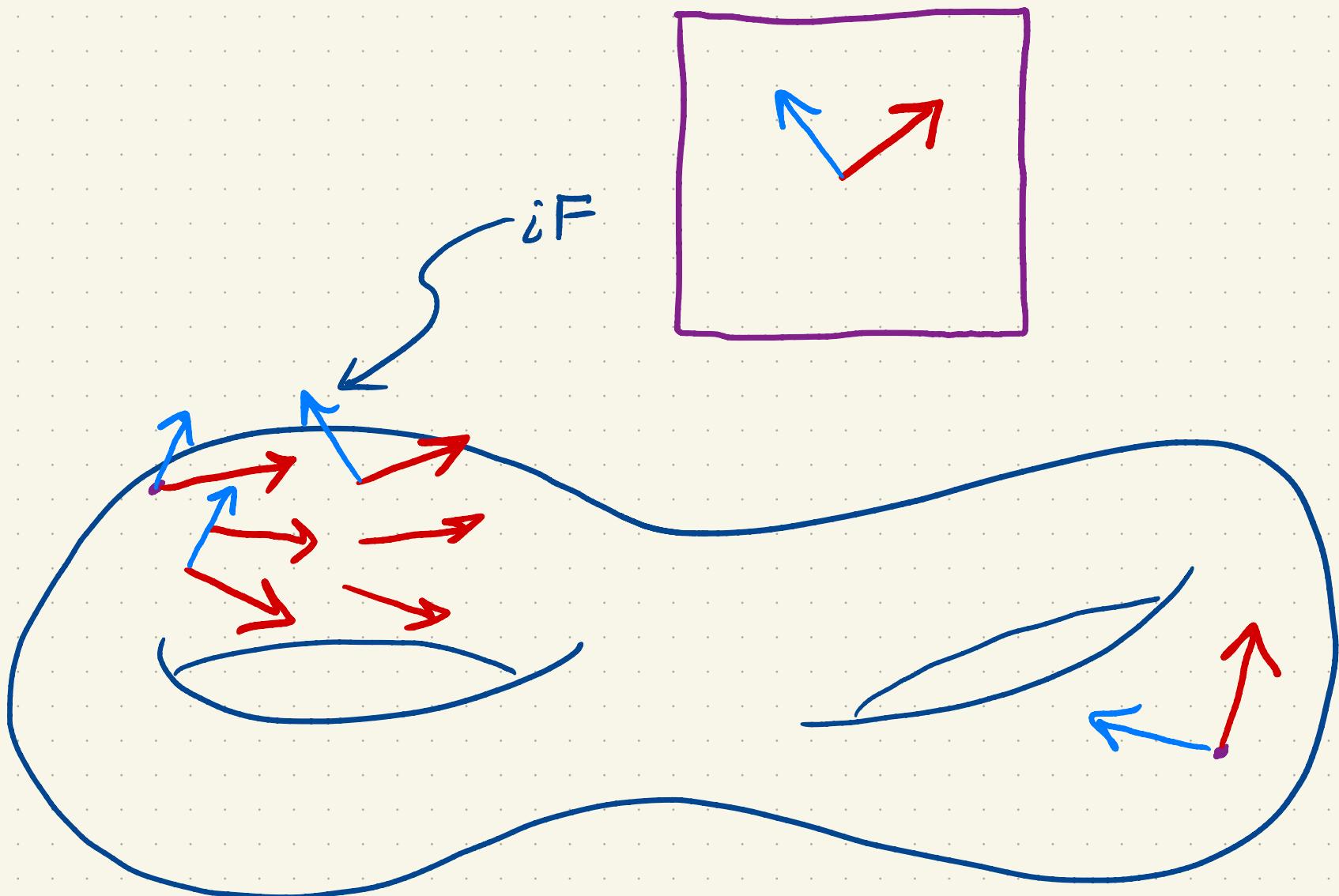
Egg Coordinates



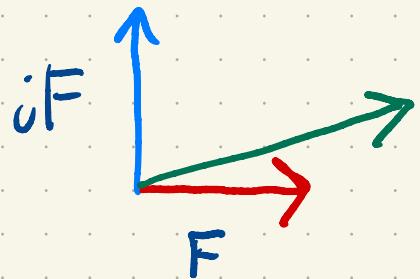
Egg Coordinates



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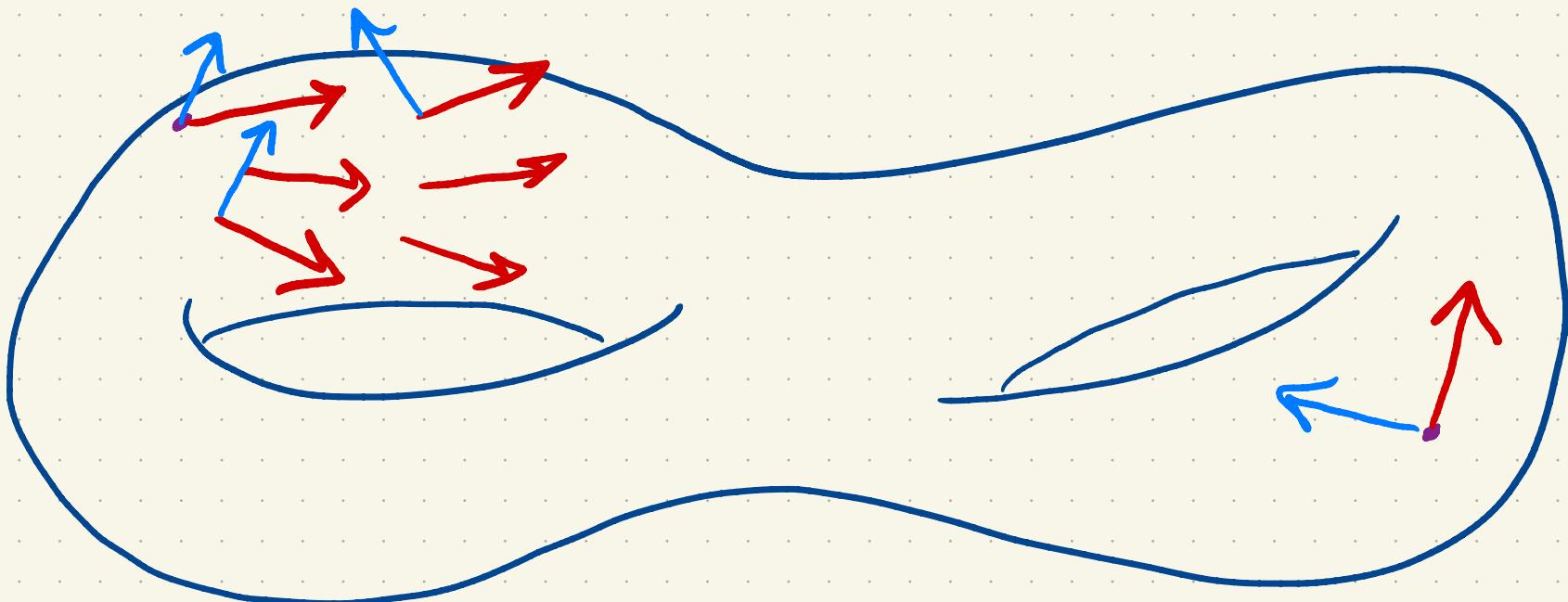
Egg Coordinates



$$Z = aF + i bF$$

$$= zF$$

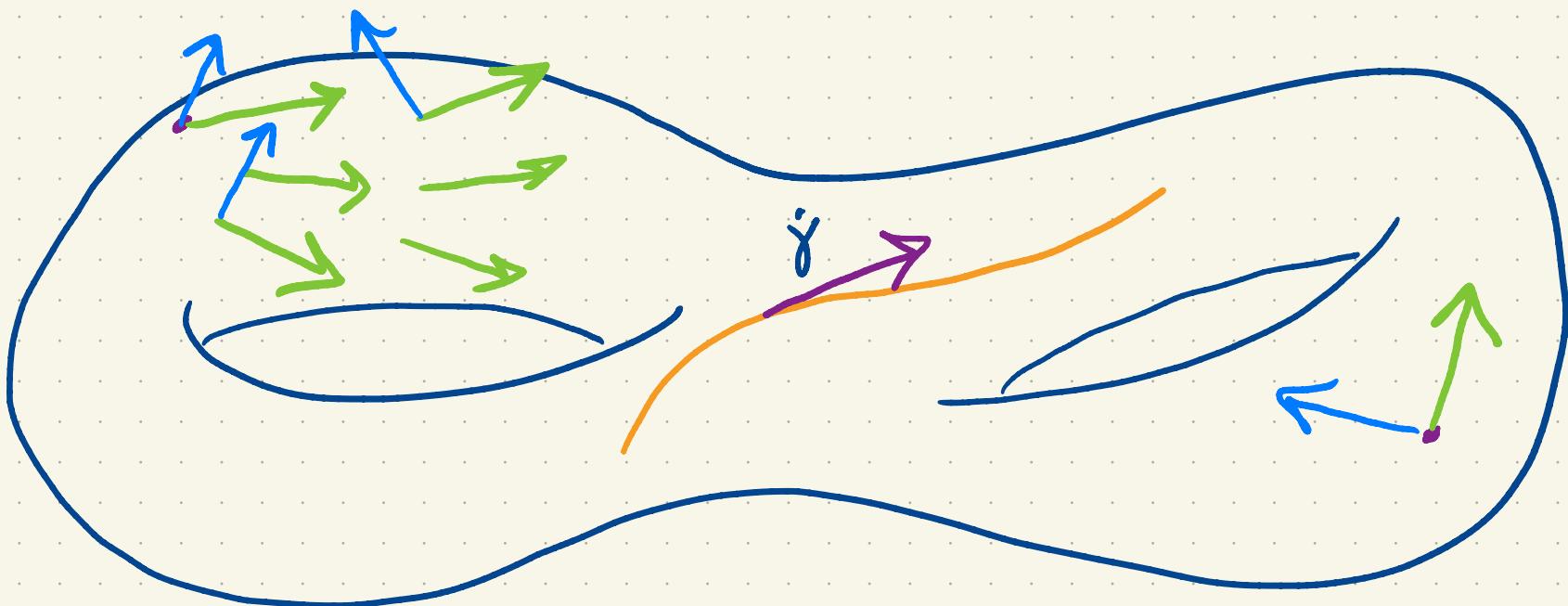
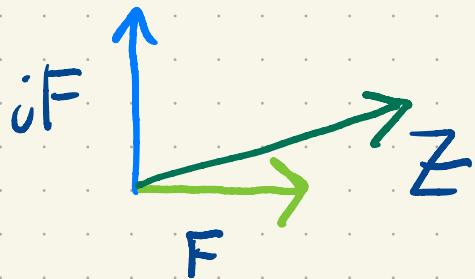
representation of
vector field via $z \in \mathbb{C}$



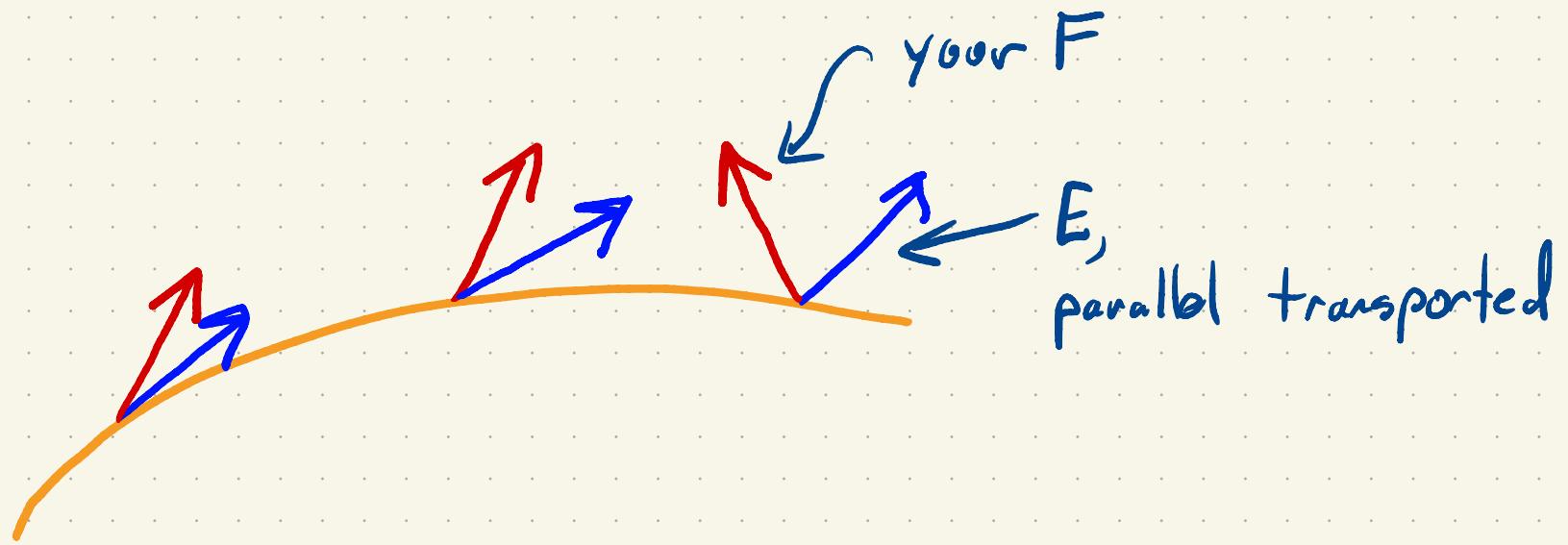
Egg Coordinates

- How to compute

$$\nabla_{\dot{y}} Z ?$$

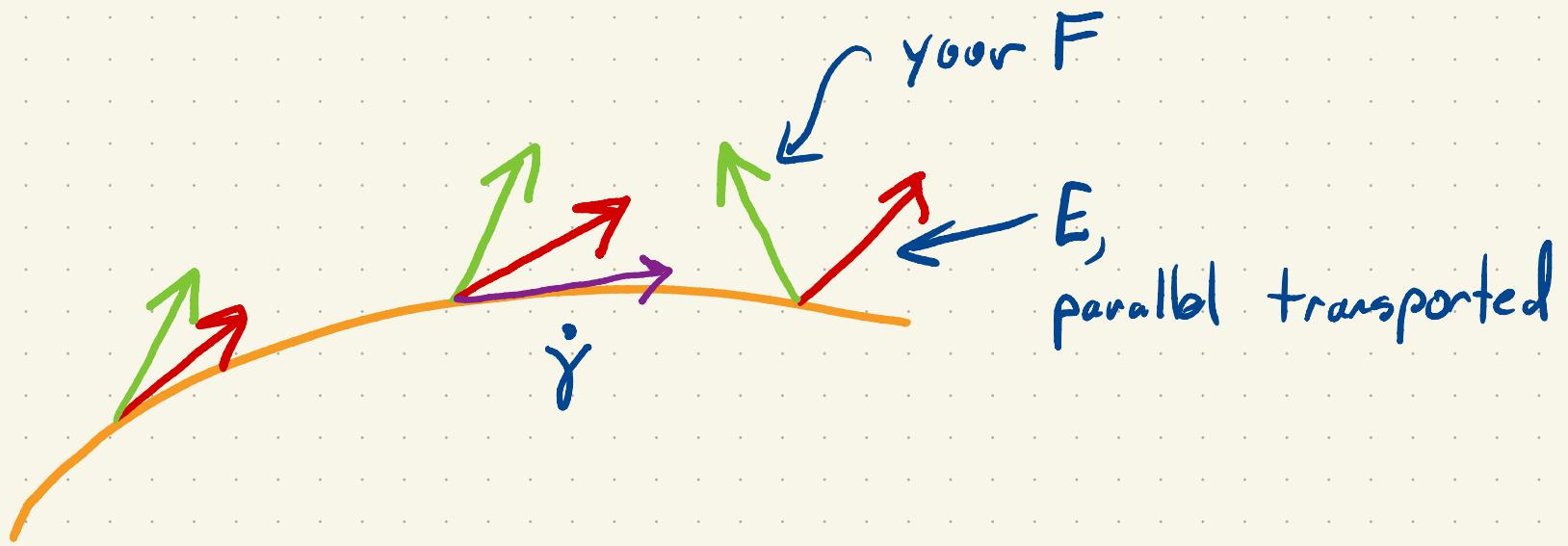


Connection 1-form



$$F = e^{i\theta} E$$

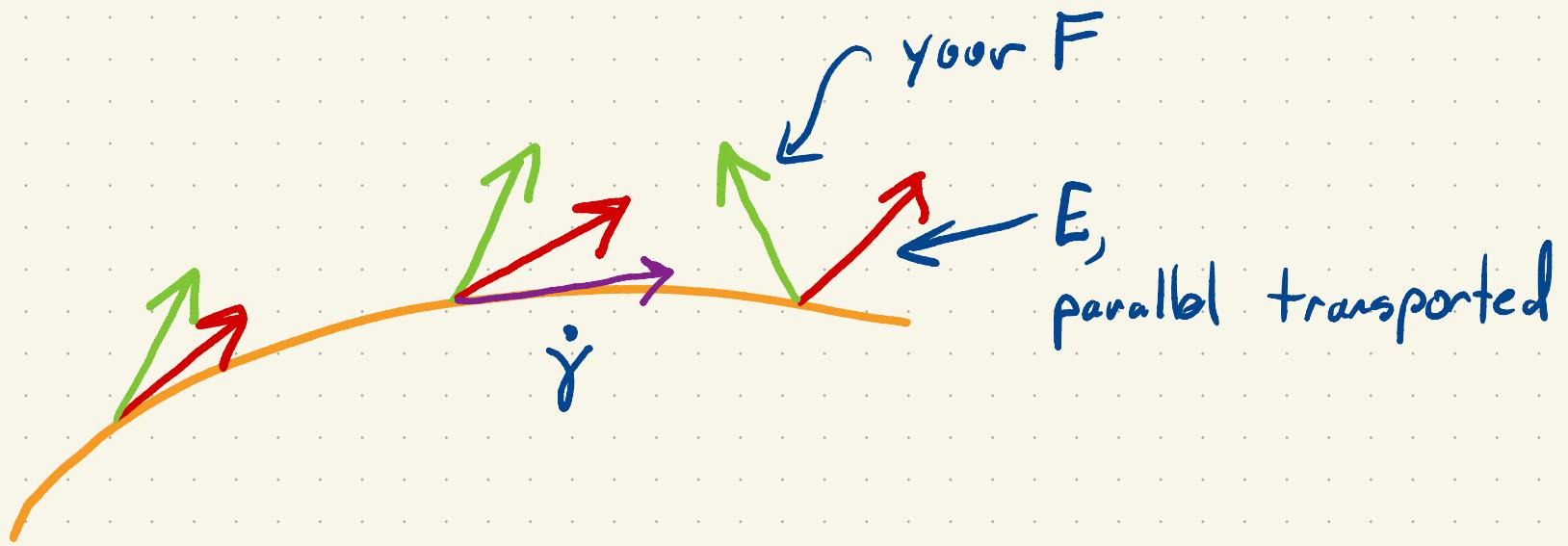
Connection 1-form



$$F = e^{i\theta} E$$

$$\nabla_{\dot{\gamma}} F = i e^{i\theta} \dot{\theta} E = i \dot{\theta} F$$

Connection 1-form

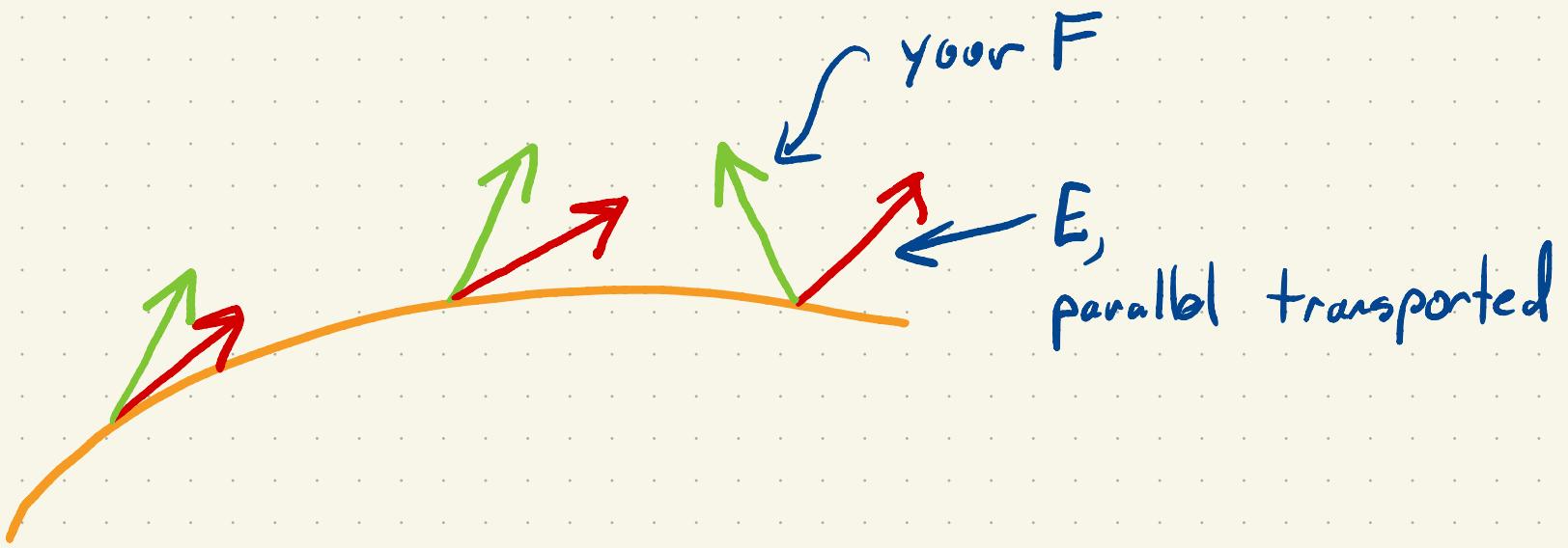


$$F = e^{i\theta} E$$

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$$A[\dot{\gamma}] = \dot{\theta}$$

Connection 1-form



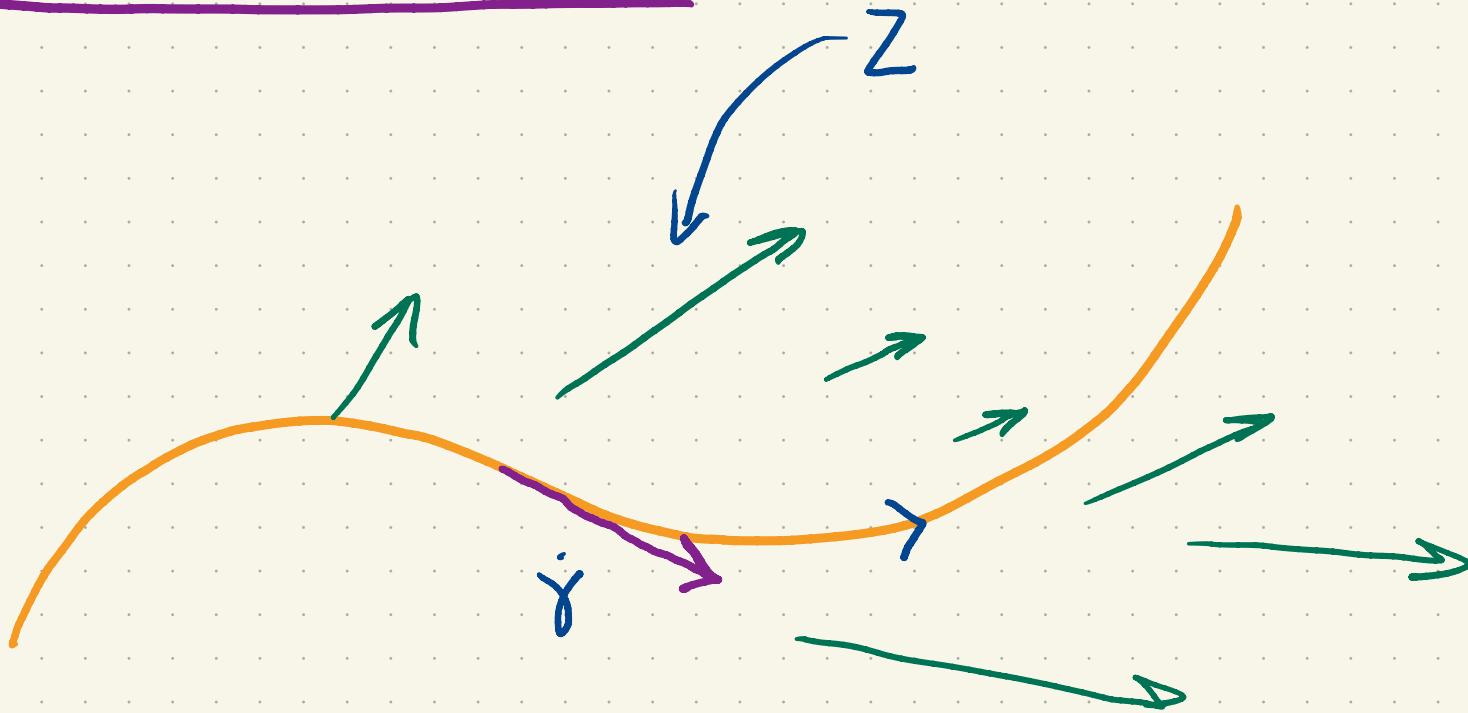
$$F = e^{i\theta} E$$

$$\nabla_{\dot{\gamma}} F = i e^{i\theta} \dot{\theta} E = i \dot{\theta} F$$

$$\nabla_X F = i A[X] F$$

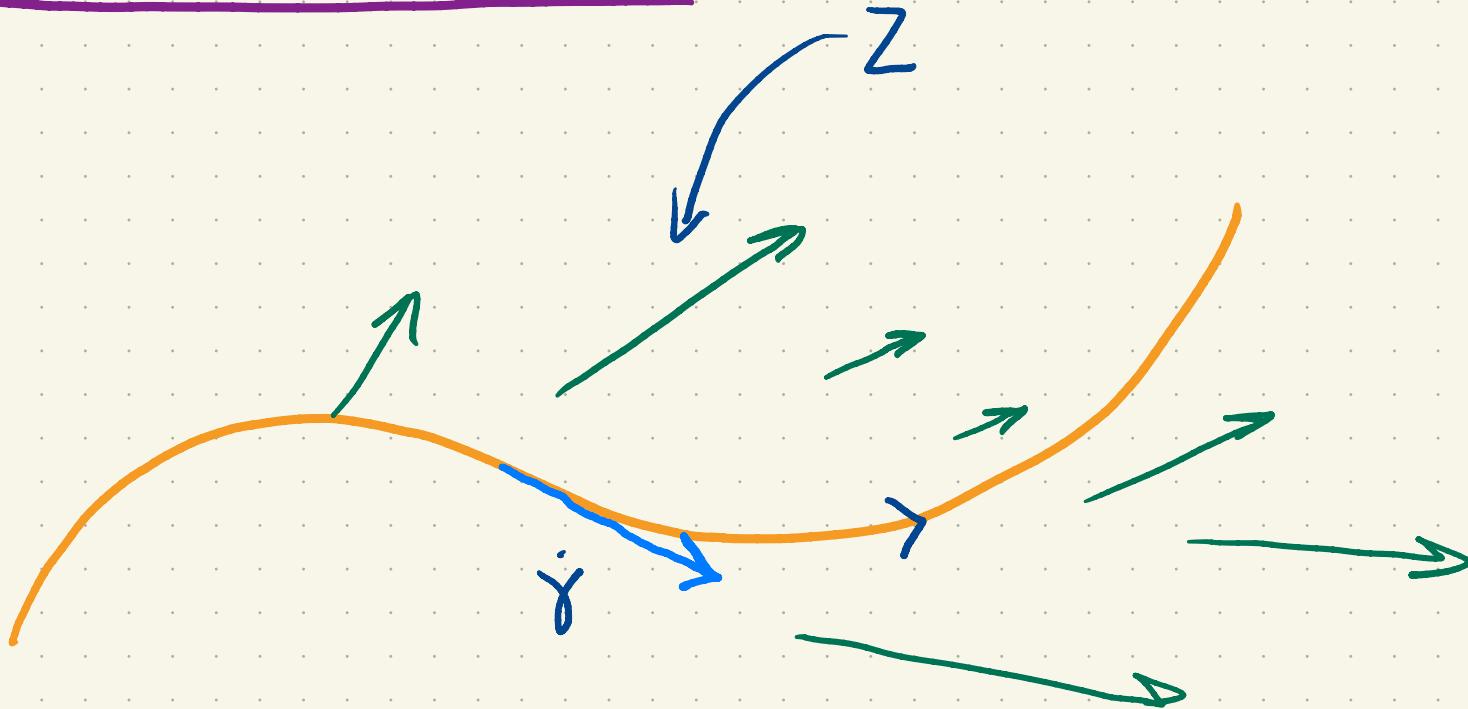
$$A[\dot{\gamma}] = \dot{\theta}$$

Derivatives Redux



$$z = z(F)$$

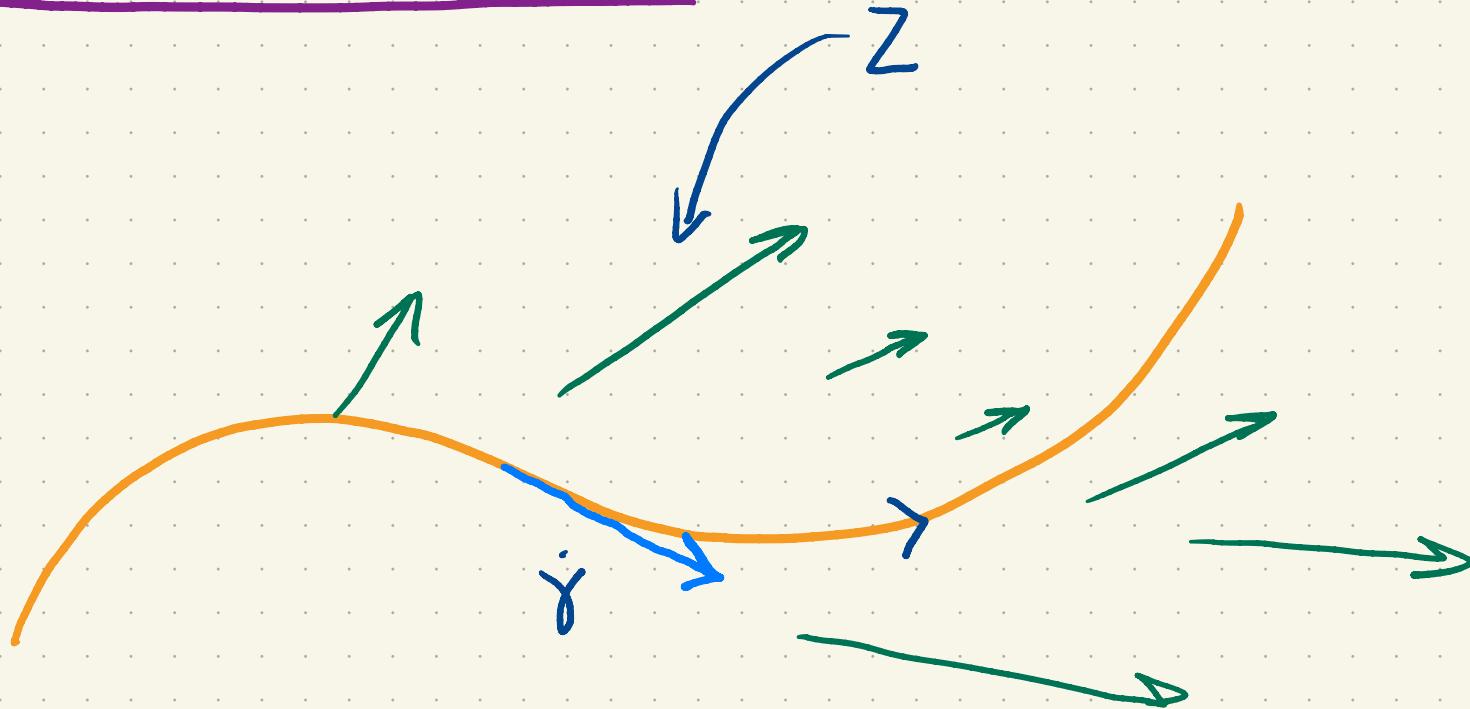
Derivatives Redux



$$z = z F$$

$$\nabla_y z = (\nabla_y z) F + z \nabla_y F$$

Derivatives Redux



$$z = z F$$

$$\nabla_{\dot{z}} z = (\nabla_{\dot{z}} z) F + z \nabla_{\dot{z}} F$$

$$= (dz[\dot{z}] + z i A[\dot{z}]) F$$

Derivative Summary

- Pick a local frame F

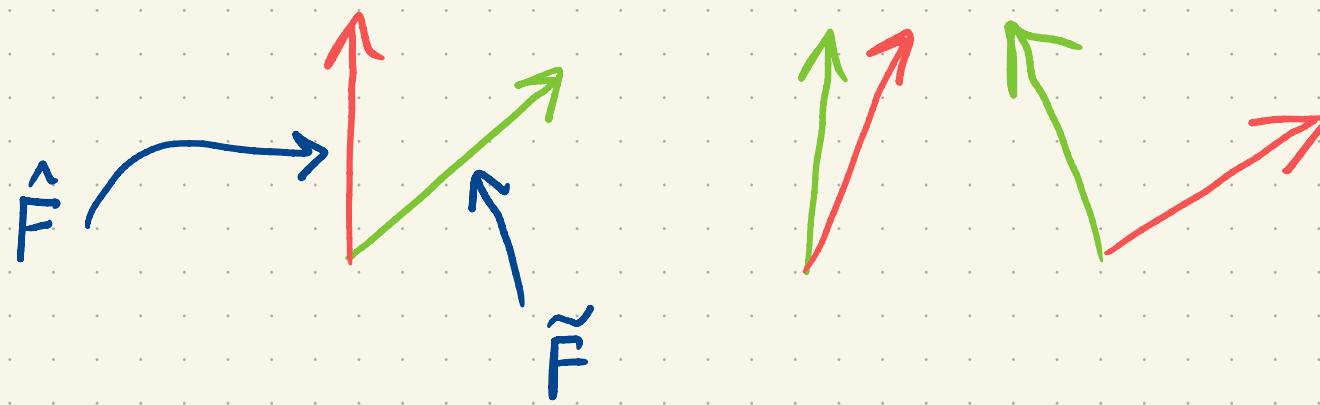
- Compute connection 1-form A
(relative to F)

$$\nabla_{\dot{\gamma}} F = i\dot{\theta}F = iA(\dot{\gamma})F$$

- Represent the vector field Z as zF

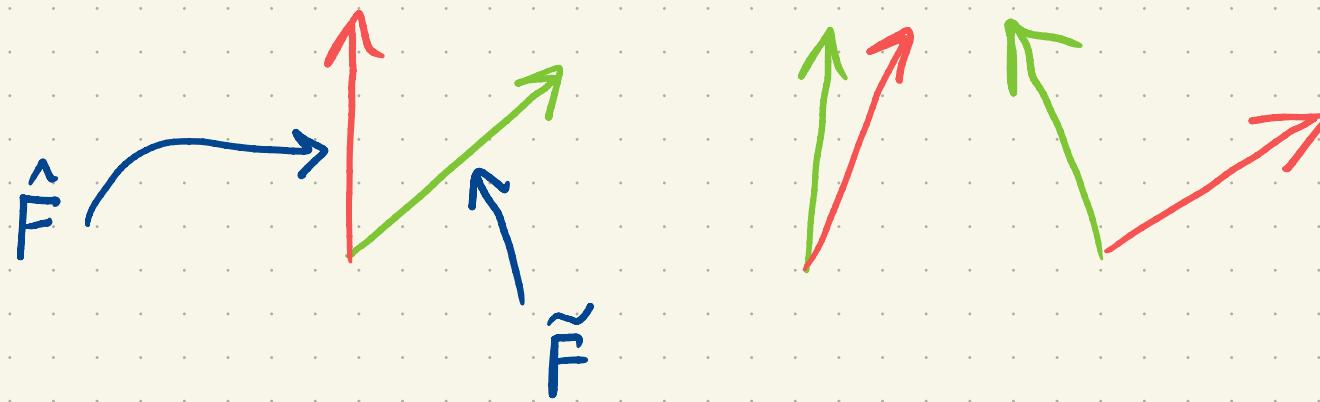
$$\nabla_X Z = [dz(X) + iA(X)]F$$

Change of Frame



$$\hat{F} = e^{i\psi} \tilde{F}$$

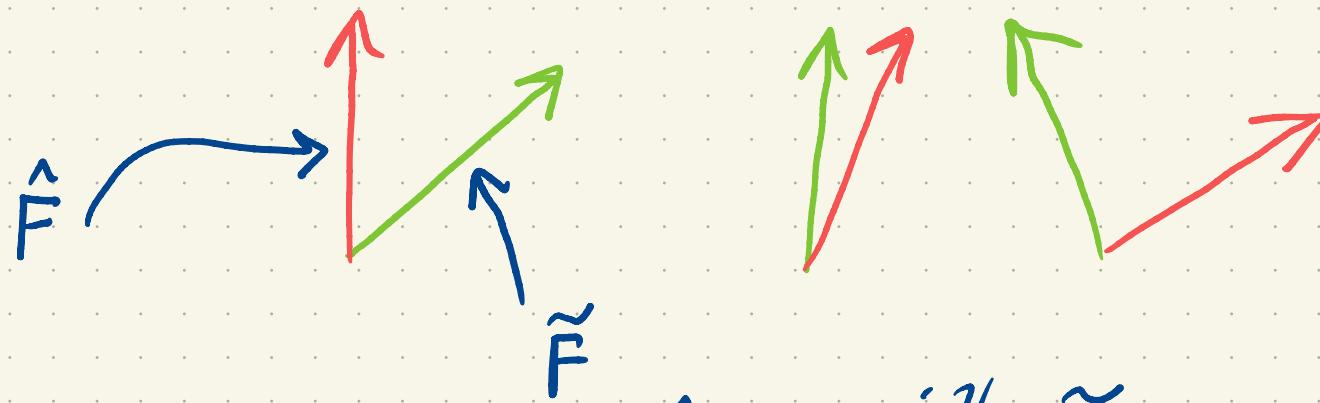
Change of Frame



$$\hat{F} = e^{i\gamma} \tilde{F}$$

$$\hat{\nabla}_X \hat{F} = i \hat{A}[X] \hat{F}$$

Change of Frame

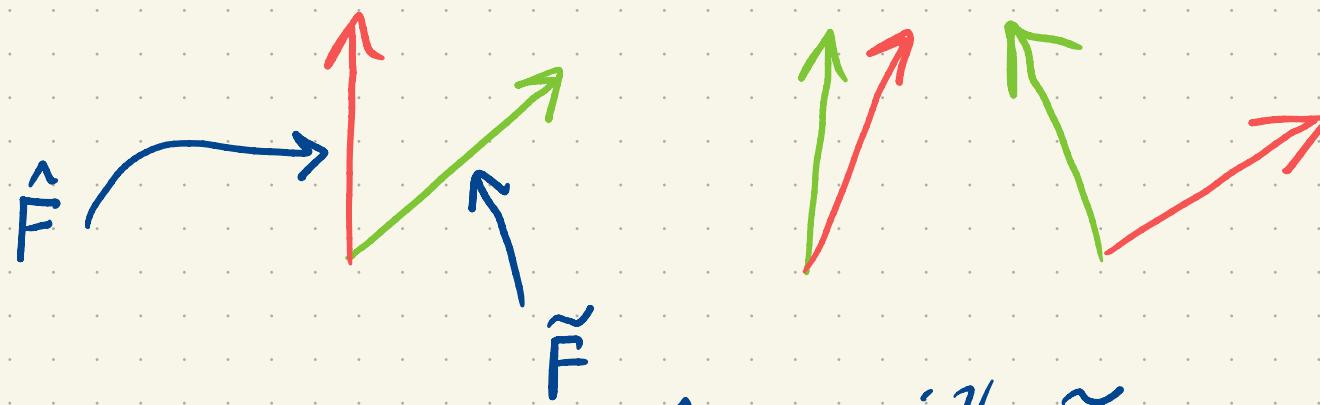


$$\hat{F} = e^{i\gamma} \tilde{F}$$

$$\nabla_X \hat{F} = i \hat{A}[X] F$$

$$\nabla_X \hat{F} = \nabla_X (e^{i\gamma} \tilde{F})$$

Change of Frame



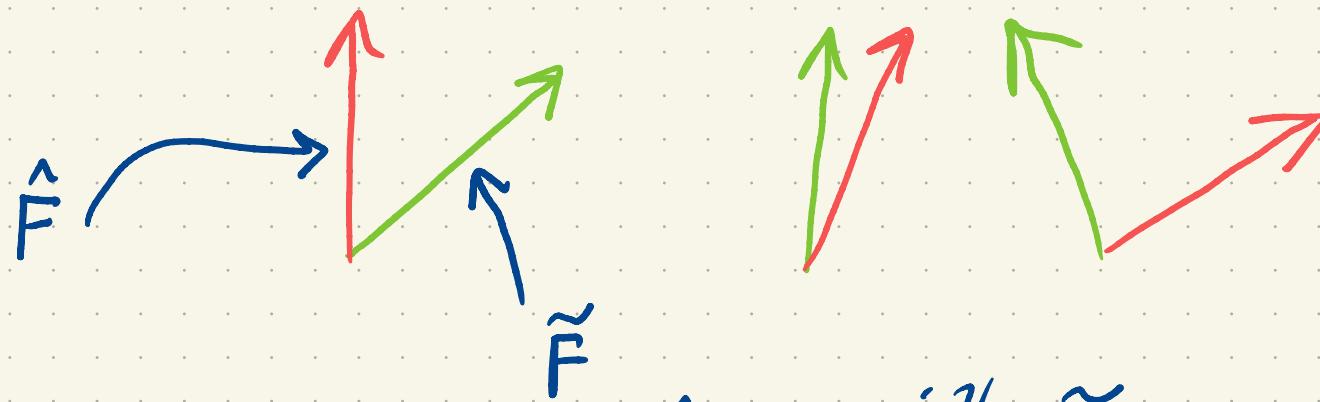
$$\hat{F} = e^{i\gamma} \tilde{F}$$

$$\nabla_X \hat{F} = i \hat{A}[X] F$$

$$\nabla_X \hat{F} = \nabla_X (e^{i\gamma} \tilde{F})$$

$$= id\gamma[X] e^{i\gamma} \tilde{F} + e^{i\gamma} i \tilde{A}[X] \tilde{F}$$

Change of Frame



$$\hat{F} = e^{i\psi} \tilde{F}$$

$$\hat{A} = d\psi + \tilde{A}$$

$$\nabla_X \hat{F} = i \hat{A}[X] F$$

$$\nabla_X \hat{F} = \nabla_X (e^{i\psi} \tilde{F})$$

$$= i d\psi[X] e^{i\psi} \tilde{F} + e^{i\psi} i \tilde{A}[X] \tilde{F}$$

$$= i (d\psi[X] + \tilde{A}[X]) \hat{F}$$

Change of Frame Summary

You	Me
Frame "Coordinates"	$\hat{F} = e^{i\varphi} \tilde{F}$
Connection 1-form	$\hat{\alpha} = d\varphi + \tilde{\alpha}$
Vector Field	$\hat{z} = e^{-i\varphi} \tilde{z}$

Very Big Picture

- Spacetime has a local $U(1)$ symmetry.

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Very Big Picture

- Spacetime has a local $U(1)$ symmetry.
- Maxwell's equations govern a connection.
- After a choice of frame, the connection can be represented by a 1-form A .
- A and $A + d\gamma$ represent the same connection
- The electric and magnetic fields are related to dA .

Homogeneous Maxwell Equations

Inertial coordinates of SR: (t, x^1, x^2, x^3) , $c=1$

$$A = \phi dt - A^i dx^i \quad (\text{implicit sum})$$

$$dA = -E^i dx^i \wedge dt - B_{ij} dx^i \wedge dx^j$$

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$$\vec{A} = (A^1, A^2, A^3)$$

$$E^i = -(\partial_i \phi + \dot{A}^i)$$

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$$\vec{B} = (B^1, B^2, B^3) \quad B^1 = B_{23}, B^2 = B_{31}, B^3 = B_{12}$$

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$$\vec{B} = \vec{\nabla} \times \vec{A}$$

Homogeneous Maxwell Equations

$$dA = -E^i dx^i \wedge dt - B_{ij} dx^i \wedge dx^j$$

$$d(A + d\varphi) = dA \quad d^2 = 0$$


Homogeneous Maxwell Equations

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$$d^2 A = 0$$

$$\begin{aligned} (\partial_j E^i - \dot{\partial}_{ij}) dx^i \wedge dx^j \wedge dt \\ - \partial_k B_{ij} dx^k \wedge dx^i \wedge dx^j = 0 \end{aligned}$$

Homogeneous Maxwell Equations

$$dA = -E^i dx^i \wedge dt - B_{ij} dx^i \wedge dx^j$$

$$d^2 A = 0$$

$$\left(\partial_j E^i - \dot{B}_{ij} \right) dx^i \wedge dx^j \wedge dt - \partial_k B_{ij} dx^k \wedge dx^i \wedge dx^j = 0$$

$$\dot{\vec{B}} = -\vec{\nabla} \times \vec{E} \quad (\vec{B}_{12} \Rightarrow \vec{B}^3)$$

$$\vec{\nabla} \cdot \vec{B} = 0$$

Homogeneous Maxwell Equations

$$A + \delta \psi$$

$$dA = -E_i dx^i \wedge dt - B_{ij} dx^i \wedge dx^j$$

$$d^2 A = 0$$

$$d \mathcal{F}_i = 0$$

$$\begin{aligned} & (\partial_j E_i - \dot{\partial}_{ij}) dx^i \wedge dx^j \wedge dt \\ & - \partial_k B_{ij} dx^k \wedge dx^i \wedge dx^j = 0 \end{aligned}$$

$$\dot{\vec{B}} = -\vec{\nabla} \times \vec{E}$$

$$\vec{\nabla} \cdot \vec{B} = 0$$

$$d \mathcal{F}_i = 0$$

~~A~~