

Last class:

$\alpha(s)$: curve in spacetime

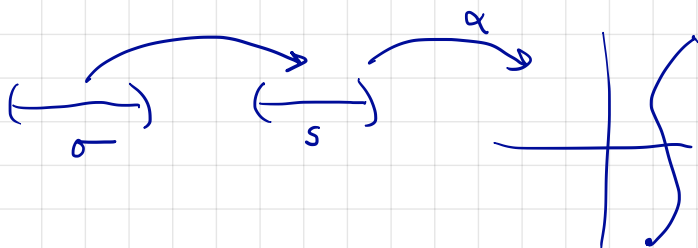
$$\alpha(s) = \begin{bmatrix} ct(s) \\ \vec{x}(s) \end{bmatrix}$$

$\alpha'(s)$ transforms as a vector. A curve is causal if α' always is.

$$\int_a^b \frac{1}{c} \sqrt{g(\alpha', \alpha')} ds = \Delta\tau, \quad \text{change in proper time along curve.}$$

elapsed time for the traveller,
from $s=a$ to $s=b$.

A curve can be reparameterized:



$$\alpha(s(\sigma))$$

We require $\frac{ds}{d\sigma} \geq 0$ so that

there is no double backing.

between points
along a causal curve

Theorem: Proper time \uparrow is independent of the parametrization

Pf: Let $\alpha(s)$ be a causal curve and consider
the proper time between $\alpha(s_0)$ and $\alpha(s_1)$.

Let

$$\beta(\sigma) = \alpha(s(\sigma)).$$

be a reparameterization with $s_0 = s(\sigma_0)$, $s_1 = s(\sigma_1)$.

$$\text{Then } \beta'(\sigma) = \alpha'(s(\sigma)) \frac{ds}{d\sigma} \text{ and}$$

$$\sqrt{g(\beta', \beta')} = \sqrt{g(\alpha', \alpha')} \frac{ds}{d\sigma}, \text{ using } \frac{ds}{d\sigma} \geq 0.$$

$$\begin{aligned} \text{By change of vars, } \int_{\sigma_0}^{\sigma_1} \frac{1}{c} \sqrt{g(\beta', \beta')} d\sigma &= \int_{\sigma_0}^{\sigma_1} \frac{1}{c} \sqrt{g(\alpha', \alpha')} \frac{ds}{d\sigma} d\sigma \\ &= \int_{s_0}^{s_1} \frac{1}{c} \sqrt{g(\alpha', \alpha')} ds. \end{aligned}$$

Def: A causal curve $\alpha(s)$ is parameterized by proper time if $g(\alpha', \alpha') = c^2$ along the curve.

Note, for such a curve,

$$\int_a^b \frac{1}{c} \sqrt{g(\alpha', \alpha')} ds = \int_a^b 1 ds = b - a.$$

How much proper time elapses from $s=a$ to $s=b$? $b-a$.

The parameter encodes this.

If α is a curve in one coord system and $\begin{bmatrix} \hat{x} \\ \hat{y} \end{bmatrix} = C^{-1} L C \begin{bmatrix} x \\ x \end{bmatrix} + Z$,

then $\hat{\alpha} = L\alpha + Z$. (The C 's are built in).

So $\hat{\alpha}' = L\alpha'$. In particular,

$$g(\hat{\alpha}', \hat{\alpha}') = g(L\alpha', L\alpha') = g(\alpha', \alpha').$$

So if α is parameterized by arclength, so is $\hat{\alpha}$.

A geometric analogy:

Def: A curve in the plane (\mathbb{R}^2) or \mathbb{R}^3 is parameterized by arc length if $|\alpha'(s)| = 1$ for all s .

e.g. $\alpha(s) = [R \cos(s/R), R \sin(s/R)]$ is

parameterized by arc length. $\alpha'(s) = [\cos(s/R), \sin(s/R)]$.

How long to go around the circle? $2\pi R$.

How far around the circle? $2\pi R$.

As s goes up by 1, so does arc length

Given any curve in the plane with $\alpha' \neq 0$

we can reparameterize it by arclength.

$$\beta(\sigma) = \alpha(s(\sigma))$$

$$\beta'(\sigma) = \alpha'(s(\sigma)) s'(\sigma)$$

$$1 = |\alpha'(s(\sigma))| \frac{ds}{d\sigma}$$

$$\text{So } \frac{ds}{d\sigma} = \frac{1}{|\alpha'(s)|} \quad (\text{uses } |\alpha'(s)| \neq 0)$$

which is an ODE to solve for s . In fact

$$|\alpha'(s)| ds = d\sigma$$

$$\text{So } \sigma = \int_{s_0}^s |\alpha'(t)| dt.$$

Now solve for $s(\sigma)$. \rightarrow possible, in principle: $\frac{d\sigma}{ds} > 0$.

But in practice, this is awful.

- 1) Solve an impossible integral
- 2) Solve an impossible algebraic eq.

Diff'ly given a timelike curve, we can always parameterize it by proper time.

Now, given any curve $\alpha(s)$, we can determine its tangent vector at any point if it were parameterized by proper time. It's just

$$c \frac{\alpha'(s)}{\sqrt{g(\alpha', \alpha')}} \quad \text{which has length } c.$$

This rescaling is the infinitesimal version of reparameterizing

Def The 4-velocity of a timelike curve $\alpha(s)$ is

$$c \frac{\alpha'(s)}{\sqrt{g(\alpha', \alpha')}}.$$

It's just the velocity of a reparameterized curve.

When α' is timelike, I'll write $|\alpha'|$ rather than $\sqrt{g(\alpha', \alpha')}$

E.g. $\alpha(t) = \begin{bmatrix} ct \\ x(t) \end{bmatrix}$, parameterized by
coordinate time.

$$\begin{aligned} |\alpha'|^2 = g(\alpha', \alpha') &= c^2 - |\dot{x}'|^2 \\ &= c^2 - v^2 \\ &= c^2 \left(1 - \left(\frac{v}{c}\right)^2\right) \\ &= c^2 \gamma^{-2} \end{aligned}$$

$$\begin{aligned} \frac{c \alpha'}{|\alpha'|} &= \frac{c \alpha'}{c \gamma^{-1}} = \gamma \alpha' = \gamma \begin{bmatrix} c \\ x'(t) \end{bmatrix} \\ &= \gamma \begin{bmatrix} c \\ \vec{v} \end{bmatrix}. \end{aligned}$$

Your text uses notation V for $\alpha(s)$'s 4-velocity.

The 4 is old-fashioned.

And although reparameterizing is hard,

computing $\frac{d\tau}{ds}$ is easy.

$$\beta(\tau) = \alpha(s(\tau))$$

$$c = |\beta'(\tau)| = |\alpha'(s(\tau))| \frac{ds}{d\tau}$$

$$\frac{d\tau}{ds} = \frac{|\alpha'(s)|}{c}$$

$$\frac{ds}{d\tau} = \frac{c}{|\alpha'|}$$

If the curve is parameterized by coordinate time t

$$\alpha' = \begin{bmatrix} c \\ \vec{v} \end{bmatrix} \quad \text{and} \quad |\alpha'|^2 = c^2 \left(1 - \frac{v^2}{c^2}\right)$$

$$|\alpha'| = \frac{c}{\gamma(v)}$$

$$\frac{d\tau}{dt} = \gamma^{-1}(\vec{v})$$

$$\frac{dt}{d\tau} = \gamma(|v|) \quad \text{expresses time dilation.}$$

E.g. Radial motion

$$x(t) = \begin{bmatrix} ct \\ R \cos(\omega t) \\ R \sin(\omega t) \\ 0 \end{bmatrix}$$

$$\omega \leq \frac{c}{R}$$

for otherwise you travel $2\pi R$

in time $\frac{2\pi}{\omega}$ with speed $R\omega$.

$$\begin{aligned} \frac{dZ}{dt} &= \frac{1}{c} \left| \frac{dx}{dt} \right| = \frac{1}{c} \sqrt{c^2 - R^2 \omega^2} \\ &= \sqrt{1 - \left(\frac{R\omega}{c}\right)^2} \end{aligned}$$

Or: $|v|^2 = (R\omega)^2$ so $\frac{dZ}{dt} = \gamma(|v|) =$

How much proper time elapses in a single rotation?

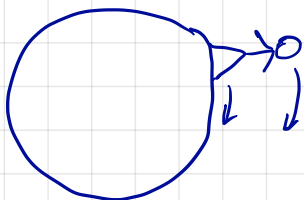
Time for a rotation: $\frac{2\pi}{\omega}$

$$\frac{2\pi}{\omega} \sqrt{1 - \left(\frac{R\omega}{c}\right)^2} \approx \frac{2\pi}{\omega} \text{ if } |R\omega| \ll c.$$

Consequence:

Orbiting at light speed, time does not pass

Consequence



Head travels faster.

Your head is younger than your feet

Next HW: how much younger?

Curvature of a curve in the plane

α'' mixes up what the curve is doing with how you parameterize it.

How curvy is it?

a) Reparam by arclength

b) $|\alpha''|$ tells your curvature.

e.g. $\alpha(s) = (R \cos(\omega s), R \sin(\omega s))$



$$\alpha'' = -\omega^2 (R \cos(\omega s), R \sin(\omega s))$$

$$|\alpha''| = \omega^2 R$$

part for the curve's radius

part for the parametrization.

We can eliminate the parametrization dependence by looking at curves parametrized by arclength.

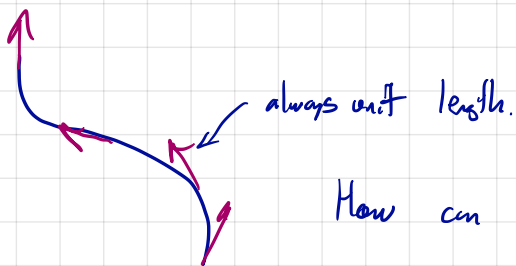
$$\omega = \frac{1}{R}$$

$$\text{Now } \alpha'' = -\frac{1}{R^2} \alpha \quad \text{and } |\alpha''| = \frac{1}{R}$$

Tiny circle, huge $\frac{1}{R}$. Big circle, tiny $\frac{1}{R}$.

We call this quantity $|\alpha''|$ when α is parametrized by arclength, the curvature of the curve.

Units: $1/L$.



always unit length.

How can α' change?

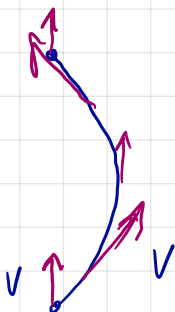
$$\alpha' \cdot \alpha' = 1$$

$$\text{so } \frac{d}{ds} \alpha' \cdot \alpha' = 0$$

$$\hookrightarrow = 2\alpha' \cdot \alpha''$$

So α'' is always perp to α' . All α' can do is rotate.

Space like picture



$$|v| = c$$

Analogously,

$$4\text{-velocity: } \frac{d\alpha}{d\tau} = V$$

$$4\text{-acceleration: } \frac{d^2\alpha}{d\tau^2}$$

$$V(t) = c \frac{\alpha'}{|\alpha'|}$$

$$\frac{d\tau}{dt} = \frac{1}{\gamma} = \frac{1}{\gamma(v)}$$

$$\frac{dV}{d\tau} = \frac{dV}{dt} \cdot \frac{dt}{d\tau} = \frac{dV}{dt} \cdot \frac{c}{|\alpha'|}$$

Typically a mess

$$\alpha = \begin{bmatrix} ct \\ R \cos(\omega t) \\ R \sin(\omega t) \end{bmatrix}$$

$$|\alpha'| = \sqrt{c^2 - R^2 \omega^2} \\ = c \sqrt{1 - \left(\frac{R\omega}{c}\right)^2}$$

$$V = \frac{1}{\sqrt{1 - \left(\frac{R\omega}{c}\right)^2}} \begin{bmatrix} c \\ -R\omega \sin \\ R\omega \cos \end{bmatrix}$$

$$\frac{dV}{d\tau} = \frac{1}{\sqrt{1 - \left(\frac{R\omega}{c}\right)^2}} \begin{bmatrix} 0 \\ -R\omega^2 \cos(\omega t) \\ -R\omega^2 \sin(\omega t) \end{bmatrix} \frac{1}{\sqrt{1 - \left(\frac{R\omega}{c}\right)^2}}$$

$$= \frac{1}{(1 - (Rv/c)^2)} \begin{bmatrix} 0 \\ -\omega^2 R \cos(\omega t) \\ -\omega^2 R \sin(\omega t) \end{bmatrix}$$

↑
relativistic correction factor

↑
usual acceleration
Newtonian

Is infinite if $R\omega = c$!

e.g.

$$\alpha(t) = \begin{bmatrix} ct \\ x(t) \end{bmatrix}$$

$$\frac{d\alpha}{dz} = \gamma(v) \begin{bmatrix} c \\ v \end{bmatrix}$$

$$\frac{dz}{dt} = \frac{1}{\gamma} \sqrt{c^2 - |v|^2}$$

$$= \gamma^{-1}$$

$$\frac{d^2 \alpha}{dz^2} = \frac{dt}{dz} \frac{d}{dt} \left[\gamma \begin{bmatrix} c \\ v \end{bmatrix} \right]$$

$$dt = \gamma dz$$

$$= \gamma \left[\gamma' \begin{bmatrix} c \\ v \end{bmatrix} + \gamma \begin{bmatrix} 0 \\ v' \end{bmatrix} \right]$$

$$= (\gamma \gamma') \begin{bmatrix} c \\ v \end{bmatrix} + \gamma^2 \begin{bmatrix} 0 \\ v' \end{bmatrix}$$

$$\frac{d}{dt} \frac{1}{\sqrt{1-\frac{v}{c}}^2} = -\frac{1}{2} \frac{1}{\left(1-\frac{v}{c}\right)^{3/2}} (-2) \frac{v}{c} \frac{v'}{c}$$
$$= \gamma^3 \frac{v}{c^2} \frac{dv}{dt}$$

$$\gamma \gamma' = \gamma^4 \frac{|v|}{c^2} \frac{d|v|}{dt}$$